

A Dynamic Problem of Thermoelasticity

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A space-periodic problem of nonlinear thermoelasticity is considered. For an elastic, linear, isotropic, homogeneous, nonviscous body in small geometry, we obtain a nonlinear system of equations. For small coefficient of the heat extension a we find a time-global weak solution of the initial-value problem. The smallness of a is independent of the length of the time interval and of the data. The space periodicity of the solution is related to the absence of reflected waves. A mixed problem for a bounded domain, even with a smooth boundary, seems to be an open problem. Our work is closely related to that by J. Nečas [5] and by J. Nečas, A. Novotný and V. Sverák [6].

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0. Introduction

Let $\Omega = (0,1)^3 \subset \mathbb{R}^3$ be the body considered. We denote:

- $u = (u_1, u_2, u_3)$ - the displacement vector,
- $\sigma = (\sigma_{ij})$ - the stress tensor,
- $e = (e_{ij})$ - the small strain tensor,
- T - the temperature,
- S - the entropy,
- $c = (c_1, c_2, c_3)$ - the heat flow,
- $f = (f_1, f_2, f_2)$ - the body force,
- W - the internal energy,
- F - the free energy,
- k - the coefficient of heat conductivity,
- d - the coefficient of specific heat,
- a - the coefficient of heat extension,
- μ, λ - the Lamé's coefficients,
- ρ - the density.

For the linear isotropic, homogeneous body, where we also take care of heat effects, we have the following equations:

$$\sigma_{ij} = \lambda \delta_{ij} e_{kk} + 2\mu e_{ij} - (3\lambda + 2\mu)a \delta_{ij}(T - T_0) \quad (\text{Hooke's law}), \quad (0.1)$$

$$\rho \ddot{u}_i = \partial \sigma_{ij} / \partial x_j + f_i \quad (\text{equation of motion}), \quad (0.2)$$

$$W = \sigma_{ij} e_{ij} - \partial c_i / \partial x_i \quad (\text{energy equation}), \quad (0.3)$$

where T_0 is a constant and δ_{ij} denotes Kronecker's delta. For simplicity we suppose $\rho = 1$. We use the summation convention over repeated indices. Let us assume that the free

energy is only a function of the temperature and of the deformation, i.e. $F(e, T) = W(e, S) - TS$. From this, by some calculations, we get

$$\dot{W} = \sigma_{ij} \dot{e}_{ij} + d\dot{T} - (3\lambda + 2\mu) a \dot{e}_{kk} T. \quad (0.4)$$

If we use the equations (0.3) and (0.4) and the relation

$$c_i(T, \nabla T) = -k(T, \nabla T) \partial T / \partial x_i \quad (0.5)$$

we obtain

$$\frac{\partial}{\partial x_i} \left(k \frac{\partial T}{\partial x_i} \right) = d\dot{T} + (3\lambda + 2\mu) a \dot{e}_{kk} T. \quad (0.6)$$

Relation (0.5) is motivated by Fourier's law. The heat flow is subject to the principle of material frame indifference and so $k(T, \nabla T)$ has to be an isotropic function (see [9]). Further we require that k is even with respect to $(T - T_0)$. The most simple function satisfying these properties is

$$k(T, \nabla T) = a_0 + a_1(T - T_0)^2 + |\nabla T|^2, \quad \text{with } a_0, a_1 = \text{const.} \quad (0.7)$$

For simplicity we set $a_0 = a_1 = 1$. If we substitute the definition of the small strain tensor into (0.1) and the result into (0.2) we obtain

$$\lambda \frac{\partial^2 u_j}{\partial x_i \partial x_j} + \mu \left(\frac{\partial^2 u_i}{\partial x_j \partial x_j} + \frac{\partial^2 u_j}{\partial x_i \partial x_j} \right) - (3\lambda + 2\mu) a \frac{\partial T}{\partial x_i} + f_i = \ddot{u}_i. \quad (0.8)$$

The system (0.6), (0.8) is a thermoelastic description of a linear, isotropic, homogeneous and thermoextensible body without viscosity. So we seek for the displacement vector u and the temperature T , both defined on $Q_I = \Omega \times I$, with $I = (0, t_0)$ and satisfying (0.6), (0.8).

By $(L^p(\Omega), \|\cdot\|_{p, \Omega})$ and $(W_p^m(\Omega), \|\cdot\|_{m, p, \Omega})$ we denote the Lebesgue and Sobolev spaces of space-periodic functions with period 1. We put

$$L^p(\Omega, \mathbb{R}^3) = L^p(\Omega) \times L^p(\Omega) \times L^p(\Omega), \quad W_p^m(\Omega, \mathbb{R}^3) = W_p^m(\Omega) \times W_p^m(\Omega) \times W_p^m(\Omega)$$

with norms defined by the same symbols $\|\cdot\|_{p, \Omega}$ and $\|\cdot\|_{m, p, \Omega}$. Let V be a Banach space. Then by $L^p(I, V)$ we denote the space of Bochner-measurable functions with values in V , for which $\int_I \|u(s)\|_V^p ds$ is finite. With the norm $\|u\|_{p, V} = \left(\int_I \|u(s)\|_V^p ds \right)^{1/p}$, $L^p(I, V)$ is a Banach space. For details see [1]. Further we denote $\dot{u} = \partial u / \partial t$, $\nabla u = (\partial u / \partial x_1, \partial u / \partial x_2, \partial u / \partial x_3)$, $\nabla^2 u = (\partial^2 u / \partial x_1^2, \partial^2 u / \partial x_1 \partial x_2, \dots, \partial^2 u / \partial x_3^2)$, weak convergence of a sequence $\{x_n\}$ to x by $x_n \rightarrow x$, $(u, v) = \int_\Omega u(x)v(x) dx$, $\langle g, h \rangle = \int g(t)h(t) dt$, $Q_t = \Omega \times (0, t)$, $t \in I$. At last, we will also denote by c constants in various estimations.

1. Existence of a solution

From the above, we have to find functions $u: Q_I \rightarrow \mathbb{R}^3$ and $T: Q_I \rightarrow \mathbb{R}$ solving the problem

$$(\lambda + \mu) \frac{\partial e_{kk}}{\partial x_i} + \mu \Delta u_i - \alpha \frac{\partial T}{\partial x_i} + f_i = \ddot{u}_i, \quad (1.1)$$

$$\beta T + \alpha e_{kk} T - \frac{\partial}{\partial x_i} \left(k \frac{\partial T}{\partial x_i} \right) = 0, \quad (1.2)$$

where $\beta > 0, \alpha \in (0, \alpha_0), \alpha_0 = \text{const}$ and, according to (0.7),

$$k = 1 + (T - T_0)^2 + |\nabla T|^2, \tag{1.3}$$

with initial conditions

$$u(x, 0) = u_0(x) \text{ and } \dot{u}(x, 0) = u_1(x), \quad x \in \Omega, \tag{1.4}$$

$$T(x, 0) = T_0 = \text{const.} > 0, \quad x \in \Omega. \tag{1.5}$$

Further we will assume that

$$u_0 \in W_2^2(\Omega, \mathbb{R}^3), \quad u_1 \in W_2^1(\Omega, \mathbb{R}^3), \quad f \in L^2(I, W_2^1(\Omega, \mathbb{R}^3)). \tag{1.6}$$

Definition 1: A solution of the given problem (1.1) - (1.6) is a pair of functions $u: Q_I \rightarrow \mathbb{R}^3$ and $T: Q_I \rightarrow \mathbb{R}$ satisfying the following conditions:

$$(i) \quad u \in L^2(I, W_2^1(\Omega, \mathbb{R}^3)), \quad \dot{u} \in L^\infty(I, L_2(\Omega, \mathbb{R}^3)), \quad \ddot{u} \in L^2(I, (W_2^1(\Omega, \mathbb{R}^3))'),$$

$$T \in L^4(I, W_4^1(\Omega)) \cap L^2(I, W_2^2(\Omega)), \quad \dot{T} \in L^{4/3}(I, W_4^1(\Omega)).$$

$$(ii) \quad (\lambda + \mu) \frac{\partial e_{kk}}{\partial x_j}(x, t) + \mu \Delta u_j(x, t) - \alpha \frac{\partial T}{\partial x_j}(x, t) + f_j(x, e) = \ddot{u}_j(x, t) \text{ a.e. in } Q_I.$$

$$(iii) \quad \int_{\Omega} \left(\beta \dot{T}(x, t) v(x) + \alpha \dot{e}_{kk}(x, t) T(x, t) v(x) + k \frac{\partial T}{\partial x_j}(x, t) \frac{\partial v}{\partial x_j}(x) \right) dx = 0 \quad \forall v \in W_4^1(\Omega), \text{ a.e. } t \in I.$$

First, let us consider the following problem:

(P1) Let $T_1 \in L^4(I, W_4^1(\Omega))$. We seek for a weak solution u of the problem (1.1), (1.4) for fixed $T = T_1$, i.e. we look for such function $u: Q_I \rightarrow \mathbb{R}^3$ that

$$(i) \quad u \in L^\infty(I, W_2^1(\Omega, \mathbb{R}^3)), \quad \dot{u} \in L^\infty(I, L_2(\Omega, \mathbb{R}^3)), \quad \ddot{u} \in L^2(I, (W_2^1(\Omega, \mathbb{R}^3))');$$

$$(ii) \quad \int_{\Omega} \left((\lambda + \mu) \frac{\partial u_i}{\partial x_j}(x, t) \frac{\partial v_j}{\partial x_j}(x) + \mu \frac{\partial u_i}{\partial x_j}(x, t) \frac{\partial v_j}{\partial x_j}(x) + \ddot{u}_i(x, t) v_i(x) \right) dx = \int_{\Omega} \left(f_i(x, t) v_i(x) - \alpha \frac{\partial T_1}{\partial x_j}(x, t) v_j(x) \right) dx \quad \forall v \in W_2^1(\Omega, \mathbb{R}^3), \text{ a.e. } t \in I.$$

Lemma 1: There exists a unique weak solution u of Problem (P1), for which the following estimate holds:

$$\int (|\dot{u}(x, t)|^2 + |\nabla u(x, t)|^2 + |u(x, t)|^2) dx \leq c \left(\|u_0\|_{1,2,\Omega}^2 + \|u_1\|_{0,2,\Omega}^2 + \|f\|_{2,Q_I}^2 + \alpha \|\nabla T_1\|_{2,Q_I}^2 \right). \tag{1.7}$$

Proof: The classical Galerkin method gives the existence of the weak solution of this linear problem (se, e.g., [8]) ■

Theorem 1: Let $T_1 \in L^2(I, W_2^2(\Omega)) \cap L^4(I, W_4^1(\Omega))$. Then for the weak solution u of the Problem (P1) we obtain

$$u \in L^\infty(I, W_2^2(\Omega, \mathbb{R}^3)), \dot{u} \in L^\infty(I, W_2^1(\Omega, \mathbb{R}^3)), \ddot{u} \in L^2(I, L^2(\Omega, \mathbb{R}^3))$$

and

$$\int (|\nabla \dot{u}(x, t)|^2 + |\nabla^2 u(x, t)|^2) dx \leq c (\|u_0\|_{2,2,\Omega}^2 + \|u_1\|_{1,2,\Omega}^2 + \|\nabla f\|_{2,Q_I}^2 + \alpha \|\nabla^2 T_1\|_{2,Q_I}^2) \text{ for a.e. } t \in I.$$

Proof: We again use the Galerkin method. Let $\{g^K\} = \{(g_1^K, g_2^K, g_3^K)\}$ be a base of the space $W_2^2(\Omega, \mathbb{R}^3)$ satisfying $\Delta g^K = -\lambda_K g^K$ for all $K \in \mathbb{N}$. We define

$$u^N(x, t) = \sum_{K=1}^N C_{NK}(t) g^K(x), \text{ where } C_{NK}(t) = (C_{NK}^1(t), C_{NK}^2(t), C_{NK}^3(t))$$

are defined by

$$\begin{aligned} & \int_{\Omega} \left((\lambda + \mu) \frac{\partial^2 u_i^N}{\partial x_j \partial x_j}(x, t) \Delta g_i^K(x) + \mu \Delta u_i^N(x, t) \Delta g_i^K(x) - \dot{u}_i^N(x, t) \Delta g_i^K(x) \right) dx \\ & = \int_{\Omega} \left(-f_i(x, t) \Delta g_i^K(x) + \alpha \frac{\partial T_1}{\partial x_i}(x, t) \Delta g_i^K(x) \right) dx \quad \forall t \in I, \end{aligned} \quad (1.8)$$

$$C_{NK}(0) = (u_0, g^K), \quad \dot{C}_{NK}(0) = (u_1, g^K).$$

Multiplying the K -equation of (1.8) by $\dot{C}_{NK}(t)$, summing up over K and integrating over $(0, t)$, we obtain

$$\begin{aligned} & \int_{\Omega} (|\nabla \dot{u}^N(x, t)|^2 + |\nabla^2 u^N(x, t)|^2) dx \\ & \leq c (\|u_0(0)\|_{2,2,\Omega}^2 + \|u_1(0)\|_{1,2,\Omega}^2 + \|\nabla f\|_{2,Q_I}^2 + \alpha \|\nabla^2 T_1\|_{2,Q_I}^2). \end{aligned} \quad (1.9)$$

From equations (1.7) and (1.9) we get that the sequence $\{u^N\}$ is bounded in $W_2^2(\Omega, \mathbb{R}^3)$ and that the sequence $\{\dot{u}^N\}$ is bounded in $W_2^1(\Omega, \mathbb{R}^3)$. Now, by a well-known technique we obtain our assertion ■

Let us for fixed $T_1 \in L^4(I, W_4^1(\Omega))$ and every $t \in I$ define the operator $A(t): W_4^1(\Omega) \rightarrow (W_4^1(\Omega))'$ by

$$(A(t)v, w) = \int_{\Omega} \left(\frac{1}{\beta} (1 + (T_1 - T_0)^2(x, t) + |\nabla v(x)|^2) \frac{\partial v}{\partial x_i}(x) \frac{\partial w}{\partial x_i}(x) \right) dx \quad \forall v, w \in W_4^1(\Omega).$$

Lemma 2: $A(t)$ is a monotonous, continuous and bounded operator (i.e. $A(t)$ maps bounded sets into bounded sets), which satisfies

$$(A(t)v - A(t)w, v - w) \geq c \int_{\Omega} (|\nabla(v - w)(x)|^2 + |\nabla(v - w)(x)|^4) dx \quad \forall v, w \in W_4^1(\Omega).$$

Proof: Put $k_i(t, x, y) = \beta^{-1} (1 + (T_1 - T_0)^2(x, t) + |y|^2) y_i$. It is obvious that k_i satisfies the Carathéodory condition and that the estimate $|k_i(t, x, y)| \leq c (g(x, t) + \sum_{j=1}^3 |y_j|^3)$ holds for some function $g(\cdot, t) \in L^{4/3}(\Omega)$. Now it is obvious (see [1: Lemma 1.6]) that $A(t)$ is a continuous and bounded operator. Further we have

$$\begin{aligned}
 & \langle A(s)v - A(s)w, v - w \rangle \\
 & \geq \frac{1}{\beta} \int_{\Omega} \left((1 + |\nabla v|^2) \frac{\partial v}{\partial x_i} - (1 + |\nabla w|^2) \frac{\partial w}{\partial x_i} \right) \frac{\partial (v - w)}{\partial x_i} dx \\
 & = \frac{1}{\beta} \int_{\Omega} \int_0^1 (1 + 3|\nabla w + t \nabla(v - w)|^2) (\nabla(v - w), \nabla(v - w)) dt dx \\
 & \geq c \left(\int_{\Omega} \int_0^1 |\nabla(v - w)|^2 |\nabla w + t \nabla(v - w)|^2 dt dx + \int_{\Omega} |\nabla(v - w)|^2 dx \right) \\
 & \geq c \int_{\Omega} (|\nabla(v - w)|^2 + |\nabla(v - w)|^4) dx \blacksquare
 \end{aligned}$$

Put $X = L^4(I, W_4^1(\Omega))$. Then $X' = L^{4/3}(I, (W_4^1(\Omega))')$. Let us define the operator $A: X \rightarrow X'$ by

$$(Av)(t) = A(t)v(t) \quad \forall v \in X$$

and the element $b \in X'$ by

$$(b(t), v) = \frac{\alpha}{\beta} \int_{\Omega} \dot{u}_i(x, t) \left(\frac{\partial T_1}{\partial x_i}(x, t)v(x) + T_1(x, t) \frac{\partial v}{\partial x_i}(x) \right) dx \quad \forall v \in W_4^1(\Omega),$$

where u is the corresponding solution of problem (P1) to $T_1 \in X$. If we use this notation we can for (1.2), (1.3), (1.5) formulate the following problem:

(P2) Find $T_2 \in X$ satisfying

$$\begin{aligned}
 (\dot{T}_2(t), v) + (A(t)T_2(t), v) &= (b(t), v) \quad \forall v \in W_4^1(\Omega), \text{ a.e. } t \in I, \\
 T_2(x, 0) &= T_0.
 \end{aligned} \tag{1.10}$$

Definition 2: The function $T_2 \in X$ is called a *weak solution of the Problem (P2)* if T_2 satisfies (1.10) and if $\dot{T}_2 \in X'$.

Theorem 2: Let $T_2 \in X$ and u be the corresponding weak solution of the Problem (P1). Then there exists a weak solution T_2 of the Problem (P2).

Proof: Let us again use the Galerkin method. Let $\{g_K\}_{K=1}^{\infty}$ be a base of $W_4^1(\Omega)$. We define

$$T_2^N(x, t) = \sum_{K=1}^N C_{NK}(t)g_K(x) \quad \forall (x, t) \in Q_I,$$

where C_{NK} are determined by

$$\begin{aligned}
 (\dot{T}_2^N(t), g_K) + (A(t)T_2^N(t), g_K) &= (b(t), g_K) \quad \forall t \in I, \\
 C_{NK}(0) &= (T_0, g_K).
 \end{aligned} \tag{1.11}$$

From this we obtain

$$\int_0^s \left((\dot{T}_2^N(t), T_2^N(t)) + (A(t)T_2^N(t), T_2^N(t)) \right) dt = \int_0^s (b(t), T_2^N(t)) dt \quad \forall s \in I. \tag{1.12}$$

For the right-hand side we have in view of (1.7)

$$\begin{aligned}
 & \left| \int_0^s \int_{\Omega} \dot{u}_i(x, t) \frac{\partial T_1}{\partial x_i}(x, t) T_2^N(x, t) dx dt \right| \\
 & \leq \int_0^{t_0} \left(\int_{\Omega} |\dot{u}_i(x, t)|^2 dx \right)^{1/2} \left(\int_{\Omega} |\nabla T_1(x, t)|^4 dx \right)^{1/4} \left(\int_{\Omega} |T_2^N(x, t)|^4 dx \right)^{1/4} dt \\
 & \leq c(c_1 + \alpha \|\nabla T_1\|_{2, Q_I}) \|\nabla T_1\|_{4, Q_I} \left(\int_0^{t_0} \left(\int_{\Omega} |T_2^N(x, t)|^4 dx \right)^{1/3} dt \right)^{3/4} \quad (1.13) \\
 & \leq c_0(c_1 + \alpha \|\nabla T_1\|_{2, Q_I}) \|\nabla T_1\|_{4, Q_I} (\|\nabla T_2^N\|_{2, Q_I} + \|T_2^N\|_{2, Q_I}),
 \end{aligned}$$

where we have used in the last line the imbedding of $W_2^1(\Omega)$ into $L^4(\Omega)$. In the same way we get

$$\begin{aligned}
 & \left| \int_0^s \int_{\Omega} \dot{u}_i(x, t) \frac{\partial T_2^N}{\partial x_i}(x, t) T_1(x, t) dx dt \right| \\
 & \leq c_0(c_1 + \alpha \|\nabla T_1\|_{2, Q_I}) \|\nabla T_2^N\|_{4, Q_I} (\|\nabla T_1\|_{2, Q_I} + \|T_1\|_{2, Q_I}). \quad (1.14)
 \end{aligned}$$

From (1.12) we obtain, in view of (1.13), (1.14) and the monotonicity of $A(t)$,

$$\begin{aligned}
 & \int_{\Omega} |T_2^N(x, s)|^2 dx + \int_{Q_I} (|\nabla T_2^N(x, t)|^2 + |\nabla T_2^N(x, t)|^4) dx dt \\
 & \leq c(\|T_2^N(x, 0)\|_{2, \Omega}^2 + (1 + \|\nabla T_1\|_{2, Q_I}^2)(\|\nabla T_1\|_{4, Q_I}^2 + \|T_1\|_{2, Q_I}^2)).
 \end{aligned}$$

But

$$\left(\int_0^{t_0} \left(\int_{\Omega} T^2(x, t) dx \right)^2 + \int_{\Omega} |\nabla T(x, t)|^4 dx dt \right)^{1/4} \quad (1.15)$$

is an equivalent norm on the space X . And so there exist a weak convergent subsequence $\{T_2^N\} \subset X$ and elements $w \in X'$, $z \in L^2(\Omega)$ such that

$$\begin{aligned}
 & T_2^N \rightharpoonup T_2 \text{ in } X, \quad AT_2^N \rightharpoonup w \text{ in } X', \\
 & T_2^N(t_0) \rightharpoonup z \text{ in } L^2(\Omega), \quad T_2^N(0) \rightharpoonup T_0 \text{ in } W_2^1(\Omega),
 \end{aligned}$$

where we used the boundedness of A and the imbedding of the set $\{T: T \in X, \dot{T} \in X'\}$ into $C(I, L^2(\Omega))$. By some calculation we obtain from (1.11)

$$\bar{T}_2 \in X', \quad \bar{T}_2 + w = b, \quad T_2(0) = T_0, \quad T_2(t_0) = z.$$

But

$$\begin{aligned}
 \overline{\lim}_{N \rightarrow \infty} \langle AT_2^N, T_2^N \rangle &= \overline{\lim}_{N \rightarrow \infty} 2^{-1} (\|T_2^N(0)\|_{2, \Omega}^2 - \|T_2^N(t_0)\|_{2, \Omega}^2) + \langle b, T_2^N \rangle \\
 &\leq 2^{-1} (\|T_2(0)\|_{2, \Omega}^2 - \|T_2(t_0)\|_{2, \Omega}^2) + \langle b, T_2 \rangle = \langle w, T_2 \rangle.
 \end{aligned}$$

This, (1.15), the weak convergence of T_2^N to T_2 in X and the monotonicity of A imply $AT_2 = w$, i.e. T_2 is a weak solution of the Problem (P2) ■

Theorem 3: Let $T_1 \in X \cap L^2(I, W_2^2(\Omega))$. Then the weak solution T_2 of the Problem (P2) belongs to the space $L^2(I, W_2^2(\Omega))$ and satisfies the estimate

$$\|\nabla^2 T_2\|_{2, Q_I}^2 \leq c_2 (c_3 + \alpha(c_5 + \alpha \|\nabla^2 T_1\|_{2, Q_I}^2) + \|\nabla T_1\|_{4, Q_I}^2 \|\nabla T_2\|_{4, Q_I}^2) \tag{1.16}$$

Proof: (i) First we prove that $T_2 \in L^2(I, W_2^2(\Omega))$. For $v \in X$ and almost every $t \in I$ we have

$$\begin{aligned} & \int_{\Omega} \left(\dot{T}_2(x, t) v(x, t) + \frac{1}{\beta} (1 + (T_1 + T_0)^2(x, t) + |\nabla T_2(x, t)|^2) \frac{\partial T_2}{\partial x_i}(x, t) \frac{\partial v}{\partial x_i}(x, t) \right) dx \\ &= - \int_{\Omega} \frac{\alpha}{\beta} T_1(x, t) \frac{\partial \dot{u}_i}{\partial x_j}(x, t) v(x, t) dx. \end{aligned} \tag{1.17}$$

We denote $\Delta_{\tau} x = \tau e^j$, $|\tau| < h$, $h > 0$ fixed, where e^j ($j = 1, 2, 3$) is a unit vector and $\{e^1, e^2, e^3\}$ is a base of \mathbb{R}^3 . Further we denote $\Delta_{\tau} w(x, t) = \tau^{-1}(w(x + \Delta_{\tau} x, t) - w(x, t))$. From (1.17) we obtain

$$\begin{aligned} & \int_{\Omega} \left\{ (\dot{T}_2(x + \Delta_{\tau} x, t) - \dot{T}_2(x, t)) v(x, t) \right. \\ & \quad + \frac{1}{\beta} \left[(1 + (T_1 - T_0)^2(x + \Delta_{\tau} x, t) + |\nabla T_2(x + \Delta_{\tau} x, t)|^2) \frac{\partial T_2}{\partial x_j}(x + \Delta_{\tau} x, t) \right. \\ & \quad \left. \left. - (1 + (T_1 - T_0)^2(x, t) + |\nabla T_2(x, t)|^2) \frac{\partial T_2}{\partial x_j}(x, t) \right] \frac{\partial v}{\partial x_j}(x, t) \right\} dx \\ &= - \int_{\Omega} \frac{\alpha}{\beta} \left[\frac{\partial \dot{u}_i}{\partial x_j}(x + \Delta_{\tau} x, t) T_1(x + \Delta_{\tau} x, t) - \frac{\partial \dot{u}_i}{\partial x_j}(x, t) T_1(x, t) \right] v(x, t) dx. \end{aligned}$$

In particular for $v(x, t) = \Delta_{\tau} T_2(x, t)$ we get

$$\begin{aligned} & \frac{1}{2} \int_{\Omega} ((\Delta_{\tau} T_2)^2(x, t) - (\Delta_{\tau} T_2)^2(x, 0)) dx + \frac{1}{\beta} \int_{Q_I} \left\{ \frac{\partial \Delta_{\tau} T_2}{\partial x_j}(x, t) \frac{\partial \Delta_{\tau} T_2}{\partial x_j}(x, t) \right. \\ & \quad + \frac{1}{2} ((T_1 - T_0)^2(x + \Delta_{\tau} x, t) - (T_1 - T_0)^2(x, t)) \frac{\partial \Delta_{\tau} T_2}{\partial x_j}(x, t) \frac{\partial \Delta_{\tau} T_2}{\partial x_j}(x, t) \\ & \quad + \frac{1}{2} (\Delta_{\tau} (T_1 - T_0))^2(x, t) \frac{\partial \Delta_{\tau} T_2}{\partial x_j}(x, t) \left(\frac{\partial T_2}{\partial x_j}(x + \Delta_{\tau} x, t) + \frac{\partial T_2}{\partial x_j}(x, t) \right) \\ & \quad + \frac{1}{2} (|\nabla T_2(x + \Delta_{\tau} x, t)|^2 + |\nabla T_2(x, t)|^2) \frac{\partial \Delta_{\tau} T_2}{\partial x_j}(x, t) \frac{\partial \Delta_{\tau} T_2}{\partial x_j}(x, t) \\ & \quad \left. + \frac{1}{2} \Delta_{\tau} |\nabla T_2(x, t)|^2 \left(\frac{\partial T_2}{\partial x_j}(x + \Delta_{\tau} x, t) + \frac{\partial T_2}{\partial x_j}(x, t) \right) \frac{\partial \Delta_{\tau} T_2}{\partial x_j}(x, t) \right\} dx dt \\ & \leq \int_{Q_I} \frac{\alpha}{\beta} \Delta_{\tau} \left(\frac{\partial \dot{u}_i}{\partial x_j}(x, t) T_1(x, t) \right) \Delta_{\tau} T_2(x, t) dx dt \\ & \leq c \left((1 + \|\nabla^2 T_1\|_{2, Q_I}) \|T_1\|_{4, Q_I} \|\Delta_{\tau} \Delta_{\tau} T_2\|_{4, Q_I} \right) \\ & \leq c \left((1 + \|\nabla^2 T_1\|_{2, Q_I}) \frac{1}{\varepsilon} \|T_1\|_{4, Q_I}^2 + \varepsilon \|\partial \Delta T_2 / \partial x_{e^j}\|_{4, Q_I}^2 \right), \end{aligned}$$

where we used

$$\int_{Q_T} \Delta_\tau u(x, t) \Delta_\tau T(x, t) dx dt = - \int_{Q_T} u(x, t) \Delta_{-\tau} \Delta_\tau T(x, t) dx dt \quad (T \in W_4^1(\Omega)).$$

$$\|\Delta_{-\tau} \Delta_\tau T\|_{4, \Omega} \leq \|\partial \Delta_\tau T / \partial x_e J\|_{4, Q_I}$$

From this we get

$$\int_{Q_I} \left\{ \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) + ((T_1 - T_0)^2(x + \Delta_\tau x, t) - (T_1 - T_0)^2(x, t)) \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) + \Delta_\tau (T_1 - T_0)(x, t) \frac{\partial \Delta_\tau T_2}{\partial x_j} \left(\frac{\partial T_2}{\partial x_j}(x + \Delta_\tau x, t) + \frac{\partial T_2}{\partial x_j}(x, t) \right) \times ((T_1 - T_0)(x + \Delta_\tau x, t) + (T_1 - T_0)(x, t)) + \left(\frac{\partial T_2}{\partial x_j}(x + \Delta_\tau x, t) + \frac{\partial T_2}{\partial x_j}(x, t) \right)^2 \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) \right\} dx dt \tag{1.18}$$

$$\leq c(1 + \|\nabla^2 T_1\|_{2, Q_I}^2 \|T_1\|_{4, Q_I}^2).$$

The third term on the left-hand side we carry to the right-hand side, where we use the estimate

$$\int_{Q_I} \Delta_\tau (T_1 - T_0)(x, t) \left(\frac{\partial T_2}{\partial x_j}(x + \Delta_\tau x, t) + \frac{\partial T_2}{\partial x_j}(x, t) \right) \times \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) ((T_1 - T_0)(x + \Delta_\tau x, t) + (T_1 - T_0)(x, t)) dx dt$$

$$\leq c \frac{1}{\varepsilon} \|\Delta_\tau T_1\|_{4, Q_I}^2 \|\nabla T_2\|_{4, Q_I}^2 + c\varepsilon \int_{Q_I} \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) \times ((T_1 - T_0)(x + \Delta_\tau x, t) + (T_1 - T_0)(x, t)) dx dt$$

for an ε such that $c\varepsilon < 1$. So from (1.18) we can obtain

$$\int_{Q_I} \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) \frac{\partial \Delta_\tau T_2}{\partial x_j}(x, t) \leq c(1 + \|\nabla^2 T_1\|_{2, Q_I}^2 \|T_1\|_{4, Q_I}^2) \tag{1.19}$$

$$\leq c(1 + \|\nabla^2 T_1\|_{2, Q_I}^2 \|T_1\|_{4, Q_I}^2 + \|\nabla T_1\|_{4, Q_I}^2 \|\nabla T_2\|_{4, Q_I}^2).$$

We will use the following fact:

Let $u \in L^p(\Omega)$, $1 < p < \infty$, let $\Delta_\tau u \in L^p(\Omega_h)$ for all $h > 0$, $|\tau| \leq h$, and let $\|\Delta_\tau u\|_{p, \Omega_h} \leq c_1 < \infty$. Then, in sense of distributions, $\|\partial u / \partial x_e J\|_{p, \Omega_h} \leq c_1$.

Then we get from (1.19) that $T_2 \in L^2(I, W_2^2(\Omega))$.

(ii) Now we want to prove the estimate (1.16). We multiply (1.2) by T_2'' , where T_2'' is some second spatial derivative, and integrate over Q_I . Using partial integration we get

$$\begin{aligned}
 & \int_{\Omega} \frac{1}{2} (T_2'(x, t_0))^2 dx - \int_{\Omega} \frac{1}{2} (T_2'(x, 0))^2 dx \\
 & + \int_{Q_I} \left(1 + (T_1 - T_0)^2(x, t) + |\nabla T_2(x, t)|^2 \right) \frac{\partial T_2'}{\partial x_j}(x, t) \frac{\partial T_2'}{\partial x_j}(x, t) dx dt \\
 & = \int_{Q_I} \left(\alpha T_1(x, t) \frac{\partial \dot{u}_i}{\partial x_j}(x, t) T_2''(x, t) \right. \\
 & \quad \left. - \left(1 + (T_1 - T_0)^2(x, t) + |\nabla T_2(x, t)|^2 \right) \frac{\partial T_2}{\partial x_j}(x, t) \frac{\partial T_2'}{\partial x_j}(x, t) \right) dx dt \\
 & = \int_{Q_I} \left(\alpha T_1(x, t) \frac{\partial \dot{u}_i}{\partial x_j}(x, t) T_2''(x, t) - 2(T_1 - T_0)(x, t) T_1'(x, t) \frac{\partial T_2}{\partial x_j}(x, t) \frac{\partial T_2'}{\partial x_j}(x, t) \right. \\
 & \quad \left. - 2 \frac{\partial T_2}{\partial x_j}(x, t) \frac{\partial T_2'}{\partial x_j}(x, t) \frac{\partial T_2}{\partial x_t}(x, t) \frac{\partial T_2'}{\partial x_t}(x, t) \right) dx dt.
 \end{aligned}$$

The last term on the right - hand side we carry to the left - hand side and so we get

$$\begin{aligned}
 & \int_{Q_I} \left(|\nabla T_2'(x, t)|^2 + (T_1 - T_0)^2(x, t) |\nabla T_2'(x, t)|^2 + |\nabla T_2(x, t)|^2 |\nabla T_2'(x, t)|^2 \right) dx dt \\
 & \leq \int_{\Omega} |\nabla T_2(x, 0)|^2 dx + \alpha \left| \int_{Q_I} T_1(x, t) T_2''(x, t) \frac{\partial \dot{u}_i}{\partial x_j}(x, t) dx dt \right| \\
 & \quad + 2 \left| \int_{Q_I} (T_1 - T_0)(x, t) \frac{\partial T_2'}{\partial x_j}(x, t) T_1'(x, t) \frac{\partial T_2}{\partial x_j}(x, t) dx dt \right| \\
 & \leq c_4 \left(c_5 + \alpha \|T_1 T_2''\|_{2, Q_I} \|\nabla \dot{u}\|_{2, Q_I} + \|(T_1 - T_0) \nabla T_2'\|_{2, Q_I} \|T_1' \nabla T_2\|_{2, Q_I} \right) \\
 & \leq c_6 \left(c_5 + \alpha \varepsilon \|T_1 T_2''\|_{2, Q_I} + \alpha \varepsilon^{-1} (c_7 + \alpha \|\nabla^2 T_1\|_{2, Q_I}) \right. \\
 & \quad \left. + \varepsilon \|(T_1 - T_0) \nabla T_2'\|_{2, Q_I}^2 + \varepsilon^{-1} \|T_1' \nabla T_2\|_{2, Q_I}^2 \right)
 \end{aligned}$$

for all $\varepsilon > 0$. From this we obtain for some ε

$$\int_{Q_I} |\nabla T_2'(x, t)|^2 dx dt \leq c_8 \left(c_5 + \alpha (c_7 + \alpha \|\nabla^2 T_1\|_{2, Q_I}^2) + \|\nabla T_1\|_{4, Q_I}^2 \|\nabla T_2\|_{4, Q_I}^2 \right) \blacksquare$$

Lemma 3: Put, for some $R_1 > 0$,

$$K_1 = \left\{ T, T \in X, \int_{Q_I} |\nabla T(x, t)|^4 dx dt \leq R_1^4, \int_{Q_I} T^4(x, t) dx dt \leq R_1^4 \right\}.$$

Let T_1, T_2 denote the same functions as in Theorem 2. Then $B: T_1 \rightarrow T_2$ is for R_1 large enough a mapping from K_1 into K_1 .

Proof: Using the proof of Theorem 2, especially (1.13) - (1.15) we have, where all the norms are L^4 -norms,

$$\begin{aligned}
& \int_{Q_I} (T_2^4(x, t) + |\nabla T_2(x, t)|^4) dx dt \\
& \leq c \int_0^{t_0} \left\{ \left(\int_{\Omega} T_2^2(x, t) dx \right)^2 + \int_{\Omega} |\nabla T_2(x, t)|^4 dx \right\} dt \\
& \leq c_9 \left(c_{10} + \alpha (c_1 + \alpha \|\nabla T_1\|) (\|\nabla T_1\| \|\nabla T_2\| + \|\nabla T_1\| \|\nabla T_2\| + \|\nabla T_2\| \|\nabla T_1\|) \right) \\
& \leq c_{11} \left(c_{10} + \alpha c_1 (\varepsilon^{-1} \|\nabla T_1\|^2 + \varepsilon \|\nabla T_2\|^2 + \varepsilon \|T_2\|^2 + \varepsilon^{-1} \|T_1\|^2) \right. \\
& \quad \left. + \alpha^2 (\varepsilon^{-1} \|\nabla T_1\|^3 + \varepsilon \|\nabla T_2\|^2 + \|\nabla T_1\|^2 + \varepsilon \|\nabla T_2\|^4 + \varepsilon \|T_2\|^4 + \|T_1\|^3) \right)
\end{aligned}$$

for all $\varepsilon > 0$. For some ε we carry all the terms containing T_2 or ∇T_2 on the left-hand side and so for $T_1 \in K_1$ we have

$$\begin{aligned}
& \int_{Q_I} (T_2^4(x, t) + |\nabla T_2(x, t)|^4) dx dt \\
& \leq c_{12} (c_{10} + \|\nabla T_1\|^2 + \|T_1\|^2 + \|\nabla T_1\|^3 + \|T_1\|^3) \leq c_{12} (c_{10} + 2R^2 + 2R^3),
\end{aligned}$$

and the last term is smaller than R_1^4 for R_1 large enough ■

Lemma 4: Put, for some $R_2 > 0$,

$$K_2 = \{T, T \in K_1, \|\nabla^2 T\|_{2, Q_I} \leq R_2, T \in X'\}.$$

Let T_1, T_2 denote the same functions as in Theorem 3. Then $B: T_1 \rightarrow T_2$, for R_2 large enough and α small enough but α independent of R_1 and R_2 , is a mapping from K_2 into K_2 .

Proof: For $T_1 \in K_2$ we have (see (1.16))

$$\begin{aligned}
\|\nabla T_2\|_{2, Q_I} & \leq c_2 (c_3 + \alpha (c_5 + \alpha \|\nabla^2 T_1\|_{2, Q_I}) + \|\nabla T_1\|_{4, Q_I} \|\nabla T_2\|_{4, Q_I}) \\
& \leq c_2 (c_3 + \alpha c_5 + \alpha^2 \|\nabla^2 T_1\|_{2, Q_I}^2 + R_1^4) \leq c_2 (c_{13} + \alpha^2 R_2^2).
\end{aligned}$$

Now we take an α such that $c_2 \alpha^2 < 1$, and so for R_2 large enough, T_2 is an element of K_2 ■

Theorem 4: Put

$$K = \{T, T \in K_2, \|\dot{T}\|_{X'} \leq R_3\}.$$

Then the mapping $B: T_1 \rightarrow T_2$ has in K a fixed point T , i.e. T and the corresponding u are a solution of (1.1) - (1.6).

Proof: We take an R_3 such that $\|\dot{T}\|_{X'} \leq R_3$ for all $T_1 \in K_2$. K is a convex, closed and bounded subset of X . So we have to show that B is a weakly continuous operator, i.e. $x_n \rightharpoonup x$ in X implies $Bx_n \rightharpoonup Bx$ in X' . Let $\{T_1^n\} \subset K$ weakly converge to T_1 . We denote $T_2^n = B(T_1^n)$. From Lemma 4 we get that $\{T_2^n\} \subset K$. So there exists a $T_3 \in K$ such that $T_2^n \rightharpoonup$

T_3 in X , $\dot{T}_2^n \rightarrow \dot{T}_3$ in $L^4(Q_I)$. From [2: Theorem 5.1] we get the strong convergence $T_2^n \rightarrow T_3$ and $T_1^n \rightarrow T_1$ in $L^4(Q_I)$. Further we have

$$\begin{aligned} & \int_{Q_I} (\dot{T}_2^n(x, t) - \dot{T}_3(x, t))(T_2^n(x, t) - T_3(x, t)) dx dt \\ & + \int_{Q_I} \frac{1}{\beta} \left[(1 + (T_1^n - T_0)^2(x, t) + |\nabla T_2^n(x, t)|^2) \frac{\partial T_2}{\partial x_i}(x, t) \right. \\ & \left. - (1 + (T_1^n - T_0)^2(x, t) + |\nabla T_3(x, t)|^2) \frac{\partial T_3}{\partial x_i}(x, t) \right] \frac{\partial T_2^n - T_3}{\partial x_i}(x, t) dx dt \\ & = - \frac{\alpha}{\beta} \int_{Q_I} \frac{\partial \dot{u}_i^n}{\partial x_i}(x, t) T_1^n(x, t) (T_2^n(x, t) - T_3(x, t)) dx dt \\ & \quad - \int_{Q_I} \dot{T}_3(x, t) (T_2^n(x, t) - T_3(x, t)) dx dt \\ & \quad - \frac{1}{\beta} \int_{Q_I} (1 + T_1^n(x, t) + |\nabla T_3(x, t)|^2) \frac{\partial T_3}{\partial x_i}(x, t) \left] \frac{\partial T_2^n - T_3}{\partial x_i}(x, t) dx dt. \end{aligned}$$

The first term on the right-hand side we can estimate by

$$c \|\nabla \dot{u}^n\|_{2, Q_I} \|T_1^n\|_{4, Q_I} \|T_2^n - T_3\|_{4, Q_I}$$

which converges to zero in view of the boundedness of $\{T_1^n\}$ and $\{\nabla \dot{u}^n\}$. The second term on the right-hand side also converges to zero, because T_3 is an element of X' . Also the third term on the right-hand side converges to zero, because $\partial T_2^n / \partial x_i \rightarrow \partial T_3 / \partial x_i$ in $L^4(Q_I)$ and $(1 + (T_1^n - T_0)^2(x, t) + |\nabla T(x, t)|^2) \partial T_3 / \partial x_i(x, t)$ converges strongly in $L^{4/3}(Q_I)$. The left-hand side we can estimate from below by

$$\begin{aligned} & \int_{\Omega} (T_2^n(x, t_0) - T_3(x, t_0))^2 dx - \int_{\Omega} (T_2^n(x, 0) - T_3(x, 0))^2 dx \\ & + \|\nabla(T_2^n - T_3)\|_{2, Q_I}^2 + \|\nabla(T_2^n - T_3)\|_{4, Q_I}^4. \end{aligned}$$

We carry the second term on the other side, but also this term converges to zero. Thus we get that $\|T_2^n - T_3\|_X \rightarrow 0$. In view of the strong convergence of T_2^n in X we can tend with $n \rightarrow \infty$ in the equation

$$\begin{aligned} & \int_{Q_I} \dot{T}_2^n(x, t) \varphi(x, t) dx dt \\ & + \frac{1}{\beta} \int_{Q_I} (1 + T_1^n(x, t)^2 + |\nabla T_2^n(x, t)|^2) \frac{\partial T_2^n}{\partial x_i}(x, t) \frac{\partial \varphi}{\partial x_i}(x, t) dx dt \\ & = - \frac{\alpha}{\beta} \int_{Q_I} \frac{\partial \dot{u}_i^n}{\partial x_i}(x, t) T_1^n(x, t) \varphi(x, t) dx dt \quad \forall \varphi \in X. \end{aligned}$$

So we get $B(T_1) = T_3$. Altogether we have the weak convergence of T_2^n to $B(T_1)$ in K , i.e. weak continuity of B . By [7: Corollary 9.3] we get our assertion ■

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