Michael T. Lock • Jeff A. Viaclovsky

# Anti-self-dual orbifolds with cyclic quotient singularities 

Received July 19, 2012 and in revised form April 20, 2013


#### Abstract

An index theorem for the anti-self-dual deformation complex on anti-self-dual orbifolds with cyclic quotient singularities is proved. We present two applications of this theorem. The first is to compute the dimension of the deformation space of the Calderbank-Singer scalar-flat Kähler toric ALE spaces. A corollary of this is that, except for the Eguchi-Hanson metric, all of these spaces admit non-toric anti-self-dual deformations, thus yielding many new examples of anti-selfdual ALE spaces. For our second application, we compute the dimension of the deformation space of the canonical Bochner-Kähler metric on any weighted projective space $\mathbb{C P}_{(r, q, p)}^{2}$ for relatively prime integers $1<r<q<p$. A corollary of this is that, while these metrics are rigid as BochnerKähler metrics, infinitely many of them admit non-trivial self-dual deformations, yielding a large class of new examples of self-dual orbifold metrics on certain weighted projective spaces.


Keywords. Anti-self-dual metrics, index theory, orbifolds

## Contents

1. Introduction . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 2805
2. The orbifold $\Gamma$-index . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 2814
3. $\Gamma$-index for cyclic group actions . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 2818
4. Non-topological terms in the $\Gamma$-index . . . . . . . . . . . . . . . . . . . . . . . . . . . 2821
5. Explicit formula for $N(q, p)$. . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 2823
6. Index on Calderbank-Singer spaces . . . . . . . . . . . . . . . . . . . . . . . . . . . 2827
7. Index on weighted projective spaces . . . . . . . . . . . . . . . . . . . . . . . . . . . 2829

References . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 2840

## 1. Introduction

If $\left(M^{4}, g\right)$ is an oriented four-dimensional Riemannian manifold, the Hodge star operator * : $\Lambda^{2} \rightarrow \Lambda^{2}$ satisfies $*^{2}=\mathrm{Id}$, and induces the decomposition of the space of 2-forms $\Lambda^{2}=\Lambda_{+}^{2} \oplus \Lambda_{-}^{2}$, where $\Lambda_{ \pm}^{2}$ are the $\pm 1$ eigenspaces of $*$. The Weyl tensor can be viewed

[^0]Mathematics Subject Classification (2010): Primary 53C25; Secondary $58 J 20$
as an operator $\mathcal{W}_{g}: \Lambda^{2} \rightarrow \Lambda^{2}$, so this decomposition enables us to decompose the Weyl tensor as $\mathcal{W}_{g}=\mathcal{W}_{g}^{+}+\mathcal{W}_{g}^{-}$, into the self-dual and anti-self-dual Weyl tensors, respectively. The metric $g$ is called anti-self-dual if $\mathcal{W}_{g}^{+}=0$, and self-dual if $\mathcal{W}_{g}^{-}=0$. Note that, by reversing orientation, a self-dual manifold is converted into an anti-selfdual manifold, and vice versa. There are now so many known examples of anti-self-dual metrics on various compact four-manifolds, that it is difficult to give a complete list here, and we refer the reader to [Via13] for a recent list of references.

The deformation theory of anti-self-dual metrics is roughly analogous to the theory of deformation of complex structures. If $(M, g)$ is an anti-self-dual four-manifold, the anti-self-dual deformation complex is given by

$$
\begin{equation*}
\Gamma\left(T^{*} M\right) \xrightarrow{\mathcal{K}_{g}} \Gamma\left(S_{0}^{2}\left(T^{*} M\right)\right) \xrightarrow{\mathcal{D}} \Gamma\left(S_{0}^{2}\left(\Lambda_{+}^{2}\right)\right), \tag{1.1}
\end{equation*}
$$

where $\mathcal{K}_{g}$ is the conformal Killing operator defined by

$$
\begin{equation*}
\left(\mathcal{K}_{g}(\omega)\right)_{i j}=\nabla_{i} \omega_{j}+\nabla_{j} \omega_{i}-\frac{1}{2}(\delta \omega) g_{i j}, \tag{1.2}
\end{equation*}
$$

with $\delta \omega=\nabla^{i} \omega_{i}, S_{0}^{2}\left(T^{*} M\right)$ denotes traceless symmetric tensors, and $\mathcal{D}=\left(\mathcal{W}^{+}\right)_{g}^{\prime}$ is the linearized self-dual Weyl curvature operator.

If $M$ is a compact manifold then there is a formula for the index depending only upon topological quantities. The analytical index is given by

$$
\begin{equation*}
\operatorname{Ind}(M, g)=\operatorname{dim}\left(H^{0}(M, g)\right)-\operatorname{dim}\left(H^{1}(M, g)\right)+\operatorname{dim}\left(H^{2}(M, g)\right) \tag{1.3}
\end{equation*}
$$

where $H^{i}(M, g)$ is the $i$ th cohomology of the complex (1.1), for $i=0,1,2$. The index is given in terms of topology via the Atiyah-Singer index theorem

$$
\begin{equation*}
\operatorname{Ind}(M, g)=\frac{1}{2}(15 \chi(M)+29 \tau(M)), \tag{1.4}
\end{equation*}
$$

where $\chi(M)$ is the Euler characteristic and $\tau(M)$ is the signature of $M$ (see [KK92]).
The cohomology groups of the complex (1.1) yield information about the local structure of the moduli space of anti-self-dual conformal classes, which we briefly recall [Ito93, KK92]. There is a map

$$
\begin{equation*}
\Psi: H^{1}(M, g) \rightarrow H^{2}(M, g), \tag{1.5}
\end{equation*}
$$

called the Kuranishi map, which is equivariant with respect to the action of $H^{0}$, and the moduli space of anti-self-dual conformal structures near $g$ is locally isomorphic to $\Psi^{-1}(0) / H^{0}$. Therefore, if $H^{2}=0$, the moduli space is locally isomorphic to $H^{1} / H^{0}$.

We will be concerned with orbifolds in dimension four with isolated singularities modeled on $\mathbb{R}^{4} / \Gamma$, where $\Gamma$ is a finite subgroup of $\operatorname{SO}(4)$ acting freely on $\mathbb{R}^{4} \backslash\{0\}$. We will say that $(M, g)$ is a Riemannian orbifold if $g$ is a smooth metric away from the singular points, and at any singular point, the metric is locally the quotient of a smooth $\Gamma$-invariant metric on the 4-ball $B^{4}$ under the orbifold group $\Gamma$.

An orbifold vector bundle is defined in terms of orbifold charts. Over a neighborhood $U_{x}$ away from the singular points, an orbifold vector bundle is defined as a vector
bundle in the usual sense, and over a neighborhood $U_{p}=\tilde{U}_{p} / \Gamma$ around the singular point $p$, where $\tilde{U}_{p}$ is a neighborhood of the origin in $\mathbb{R}^{4}$, it is identified with the quotient of a smooth $\Gamma$-equivariant vector bundle over $\tilde{U}_{p}$. On overlaps the obvious compatibility conditions are satisfied. A smooth section of an orbifold vector bundle is a globally defined section so that on a neighborhood $U_{x}$ away from the singular points it is smooth in the ordinary sense, and on a neighborhood $U_{p}$ of a singular point, it is identified with a smooth $\Gamma$-equivariant section of the corresponding $\Gamma$-equivariant bundle over $\tilde{U}_{p}$ defining the orbifold vector bundle in that neighborhood.

Consider a Riemannian orbifold $(M, g)$. The Riemannian orbifold metric $g$ is defined to be a smooth section of the orbifold bundle $S^{2}\left(T^{*} M\right)$ that is positive definite everywhere. We denote by $\mathcal{R}_{\left(\left\{p_{i}\right\},\left\{\Gamma_{i}\right\}\right)}$ the space of all Riemannian orbifold metrics on $M$ having singular points $\left\{p_{i}\right\}_{1 \leq i \leq k}$ with corresponding orbifold groups $\left\{\Gamma_{i}\right\}_{1 \leq i \leq k}$. This space is invariant under pullbacks by diffeomorphisms and under conformal changes because the $\Gamma_{i}$-equivariance of a section over $\tilde{U}_{p_{i}}$ is preserved by these operations.

The anti-self-dual condition, $\mathcal{W}_{g}^{+}=0$, is invariant under the action of $\operatorname{Diff}^{+}(M)$, the group of orientation preserving diffeomorphisms of $M$, as well as conformally invariant. Therefore, we want to study the space of equivalence classes of anti-self-dual metrics under the action of $\operatorname{Diff}^{+}(M) \times C^{\infty}(M)$. Now, we make the following definition:
Definition 1.1. Let $M$ be an orbifold with singular points $\left\{p_{i}\right\}_{1 \leq i \leq k}$ and corresponding orbifold groups $\left\{\Gamma_{i}\right\}_{1 \leq i \leq k}$. The moduli space of anti-self-dual orbifold conformal structures is defined to be

$$
\begin{equation*}
\mathcal{M}=\left\{g \in \mathcal{R}_{\left(\left\{p_{i}\right\},\left\{\Gamma_{i}\right\}\right)}: \mathcal{W}_{g}^{+}=0\right\} / \operatorname{Diff}^{+}(M) \times C^{\infty}(M), \tag{1.6}
\end{equation*}
$$

where $g_{1}$ is equivalent to $g_{2}$ if there exists a $(\phi, f) \in \operatorname{Diff}^{+}(M) \times C^{\infty}(M)$ so that $g_{1}=e^{f} \phi^{*}\left(g_{2}\right)$.
In this paper, we study the moduli space of anti-self-dual orbifold conformal structures, which from now on will be referred to only as the moduli space of anti-self-dual conformal structures.

If $(M, g)$ is an anti-self-dual orbifold, the anti-self-dual deformation complex is the same as complex (1.1) except that vector bundles and smooth sections now refer to their orbifold counterparts. The moduli space of anti-self-dual conformal structures is again given by the zero set of the Kuranishi map $\Psi: H^{1} \rightarrow H^{2}$ modulo the action of $H^{0}$, and the proofs in [Ito93, KK92] remain valid in the orbifold context. However, the index formula (1.4) does not hold without adding a correction term. In [Kaw81], Kawasaki proved a version of the Atiyah-Singer index theorem for orbifolds, and gave a general formula for the correction term. Our first result is an explicit formula for this correction term for the complex (1.1) in the case that $\Gamma$ is an action of a cyclic group. In order to state this, we first make some definitions.

For $1 \leq q<p$ relatively prime integers, we denote by $\Gamma_{(q, p)}$ the cyclic action

$$
\left(\begin{array}{cc}
\exp (2 \pi i k / p) & 0  \tag{1.7}\\
0 & \exp (2 \pi i k q / p)
\end{array}\right), \quad 0 \leq k<p
$$

acting on $\mathbb{R}^{4}$, which we identify with $\mathbb{C}^{2}$ using $z_{1}=x_{1}+i x_{2}$ and $z_{2}=x_{3}+i x_{4}$.

We will also refer to this action as a $(q, p)$-action.
Definition 1.2. A group action $\Gamma_{1}: G \rightarrow \mathrm{SO}(4)$ is conjugate to another group action $\Gamma_{2}: G \rightarrow \mathrm{SO}(4)$ if there exists an element $O \in O(4)$ such that for any $g \in G$, we have $\Gamma_{1}(g) \circ O=O \circ \Gamma_{2}(g)$. If $O \in \mathrm{SO}(4)$, then the actions are said to be orientationpreserving conjugate, while if $O \notin \mathrm{SO}(4)$, the actions are orientation-reversing conjugate.

Remark 1.3. We note the important fact that if $\Gamma$ is an $\mathrm{SO}(4)$ representation of a cyclic group, then $\Gamma$ is orientation-preserving conjugate to a $\Gamma_{(q, p)}$-action [McC02]; we therefore only need to consider the $\Gamma_{(q, p)}$-actions. Furthermore, for $1 \leq q, q^{\prime}<p$, if a $\Gamma_{(q, p)}{ }^{-}$ action is orientation-preserving conjugate to a $\Gamma_{\left(q^{\prime}, p\right)}$-action then $q q^{\prime} \equiv 1 \bmod p$. We also note that a $\Gamma_{(q, p)}$-action is orientation-reversing conjugate to a $\Gamma_{(p-q, p)}$-action.

We will employ the following modified Euclidean algorithm. For $1 \leq q<p$ relatively prime integers, write

$$
\begin{align*}
p & =e_{1} q-a_{1}, \\
q & =e_{2} a_{1}-a_{2}, \\
& \vdots  \tag{1.8}\\
a_{k-3} & =e_{k-1} a_{k-2}-1, \\
a_{k-2} & =e_{k} a_{k-1}=e_{k},
\end{align*}
$$

where $e_{i} \geq 2$, and $0 \leq a_{i}<a_{i-1}, i=1, \ldots, k$. This can also be written as the continued fraction expansion

$$
\begin{equation*}
\frac{q}{p}=\frac{1}{e_{1}-\frac{1}{e_{2}-\cdots \frac{1}{e_{k}}}} \tag{1.9}
\end{equation*}
$$

We refer to the integer $k$ as the length of the modified Euclidean algorithm.
Our main theorem expresses the correction term in the index theorem in terms of the $e_{i}$ and the length of the modified Euclidean algorithm:

Theorem 1.4. Let $(M, g)$ be a compact anti-self-dual orbifold with a single orbifold point of type $(q, p)$. The index of the anti-self-dual deformation complex on $(M, g)$ is given by

$$
\operatorname{Ind}(M, g)= \begin{cases}\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)+\sum_{i=1}^{k} 4 e_{i}-12 k-2 & \text { when } q \neq p-1  \tag{1.10}\\ \frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-4 p+4 & \text { when } q=p-1\end{cases}
$$

In some other special cases, the correction term may be written directly in terms of $p$. For example, if $q=1$ and $p>2$, we have

$$
\begin{equation*}
\sum_{i=1}^{k} 4 e_{i}-12 k-2=4 p-14 \tag{1.11}
\end{equation*}
$$

We note that the cases $q=1$ and $q=p-1$ were proved earlier in [Via13] using a different method.

Remark 1.5. While Theorem 1.4 is stated in the case of a single orbifold point for simplicity, if a compact anti-self-dual orbifold has several cyclic quotient orbifold points, then a similar formula holds, with the correction term simply being the sum of the corresponding correction terms for each type of orbifold point.

### 1.1. Asymptotically locally Euclidean spaces

Many interesting examples of anti-self-dual metrics are complete and non-compact. Given a compact Riemannian orbifold ( $\hat{X}, \hat{g}$ ) with non-negative scalar curvature, letting $G_{p}$ denote the Green's function for the conformal Laplacian associated with any point $p$, the non-compact space $X=\hat{X} \backslash\{p\}$ with metric $g_{p}=G_{p}^{2} \hat{g}$ is a complete scalar-flat orbifold. Inverted normal coordinates in the metric $\hat{g}$ in a neighborhood of $p$ give rise to a coordinate system in a neighborhood of infinity of $X$, which motivates the following:

Definition 1.6. A non-compact Riemannian orbifold ( $X^{4}, g$ ) is called asymptotically locally Euclidean or $A L E$ of order $\tau$ if there exists a finite subgroup $\Gamma \subset \mathrm{SO}(4)$ acting freely on $\mathbb{R}^{4} \backslash\{0\}$, and a diffeomorphism $\phi: X \backslash K \rightarrow\left(\mathbb{R}^{4} \backslash B(0, R)\right) / \Gamma$ where $K$ is a compact subset of $X$, satisfying $\left(\phi_{*} g\right)_{i j}=\delta_{i j}+O\left(r^{-\tau}\right)$ and $\partial^{|k|}\left(\phi_{*} g\right)_{i j}=O\left(r^{-\tau-k}\right)$ for any partial derivative of order $k$, as $r \rightarrow \infty$, where $r$ is the distance to some fixed basepoint.

An orbifold compactification of an ALE space ( $X, g$ ) is a choice of a conformal factor $u: X \rightarrow \mathbb{R}_{+}$such that $u=O\left(r^{-2}\right)$ as $r \rightarrow \infty$. This type of compactification is very important in modern geometry: see for example [Sch84, Kro89]. The regularity of the compactification is an important issue. In general, such a compactification is only a $C^{1, \alpha}$ orbifold. However, if $(X, g)$ is anti-self-dual then there moreover exists a $C^{\infty}$-orbifold conformal compactification [TV05, CLW08]. This was generalized to Bach-flat metrics in [Str10], and to obstruction-flat metrics in [AV12].

Remark 1.7. It is crucial to note that if $(X, g)$ is an anti-self-dual ALE space with a $\Gamma$-action at infinity, then the conformal compactification $(\hat{X}, \hat{g})$ with the anti-self-dual orientation has a $\tilde{\Gamma}$-action at the orbifold point where $\tilde{\Gamma}$ is orientation-reversing conjugate to $\Gamma$. In the case of a cyclic group, if the action at infinity of the anti-self-dual ALE space $(X, g)$ is of type $(q, p)$, then the action at the orbifold point of the compactification $(\hat{X}, \hat{g})$ with the anti-self-dual orientation is of type ( $p-q, p$ ).

Many examples of anti-self-dual ALE spaces with non-trivial group at infinity have been discovered. The first non-trivial example was due to Eguchi and Hanson [EH79], who found a Ricci-flat anti-self-dual metric on $\mathcal{O}(-2)$ which is ALE with group $\mathbb{Z} / 2 \mathbb{Z}$ at infinity. Gibbons-Hawking [GH78, Hit79] then wrote down a metric ansatz depending on the choice of $n$ monopole points in $\mathbb{R}^{3}$. This gives a family of anti-self-dual ALE hyperkähler metrics with cyclic actions at infinity contained in $\mathrm{SU}(2)$, which are called multi-Eguchi-Hanson metrics.

Using the Joyce construction from [Joy95], Calderbank and Singer [CS04] produced many examples of toric ALE anti-self-dual metrics, which are moreover scalarflat Kähler, and have cyclic groups at infinity contained in $\mathrm{U}(2)$. For a $(q, p)$-action, the space $X$ is the minimal Hirzebruch-Jung resolution of $\mathbb{C}^{2} / \Gamma_{(q, p)}$, with exceptional divisor given by the union of 2-spheres $S_{1} \cup \cdots \cup S_{k}$, with intersection matrix

$$
\left(S_{i} \cdot S_{j}\right)=\left(\begin{array}{ccccc}
-e_{1} & 1 & 0 & \cdots & 0  \tag{1.12}\\
1 & -e_{2} & 1 & \cdots & 0 \\
0 & 1 & -e_{3} & \cdots & 0 \\
\vdots & \vdots & \vdots & & \vdots \\
0 & 0 & 0 & \cdots & -e_{k}
\end{array}\right) \text {, }
$$

where the $e_{i}$ and $k$ are defined above in (1.8) with $e_{i} \geq 2$. The Kähler scalar-flat metric on $X$ is then written down explicitly using the Joyce ansatz from [Joy95]. We do not require the details of the construction here, but only note the following: For $q>1$ the identity component of the isometry group of these metrics is a real 2-torus, and for $q=1$, it is $U(2)$.

When $q=p-1$, these metrics are the toric Gibbons-Hawking multi-Eguchi-Hanson metrics (when all monopole points are on a common line). In this case $k=p-1$ and $e_{i}=2$ for $1 \leq i \leq k$. The moduli space of toric metrics in this case is of dimension $p-2$. But the moduli space of all multi-Eguchi-Hanson metrics is of dimension 3( $p-2$ ). So it is well known that these metrics admit non-toric anti-self-dual deformations. When $q=1$, these metrics agree with the LeBrun [LeB88] negative mass metrics on $\mathcal{O}(-p)$. In this case $k=1$ and $e_{1}=p$. For $p>2$, it was recently shown [Hon13, Via13] that these spaces also admit non-toric anti-self-dual deformations. Theorem 1.9 will give a vast generalization of this phenomenon to the general case $1<q<p-1$. The proof of Theorem 1.9 relies on the following explicit formula for the index of the complex (1.1) on the conformal compactification of these metrics:
Theorem 1.8. Let $(\hat{X}, \hat{g})$ be the orbifold conformal compactification of a CalderbankSinger space $(X, g)$ with a $(q, p)$-action at infinity. Then the index of the anti-self-dual deformation complex is given by

$$
\operatorname{Ind}(\hat{X}, \hat{g})= \begin{cases}5 k+5-\sum_{i=1}^{k} 4 e_{i} & \text { when } q \neq 1  \tag{1.13}\\ -4 p+12 & \text { when } q=1\end{cases}
$$

where $k$ and $e_{i}, 1 \leq i \leq k$, are the integers occurring in the modified Euclidean algorithm defined in (1.8).
In some other special cases we can write the index formula (1.13) directly in terms of $p$. For example, if $1<q=p-1$ the index simplifies to $-3 p+8$.

By a result of LeBrun-Maskit [LM08, Theorem 4.2], $H^{2}(\hat{X}, \hat{g})=0$ for these metrics, so the actual moduli space is locally isomorphic to $H^{1} / H^{0}$. Therefore, a direct consequence of Theorem 1.8 is that the Calderbank-Singer spaces admit large families of non-toric anti-self-dual deformations, yielding many new examples:

Theorem 1.9. Let $(X, g)$ be a Calderbank-Singer space with a $(q, p)$-action at infinity, and $(\hat{X}, \hat{g})$ be the orbifold conformal compactification. Then there exists a neighborhood, $\mathcal{U}_{\hat{g}}$, of $\hat{g}$ in the moduli space of anti-self-dual conformal structures on $\hat{X}$ such that:

- If $q=1$ and $p=2$, then $\hat{g}$ is rigid, that is, $\operatorname{dim}\left(\mathcal{U}_{\hat{g}}\right)=0$.
- If $q=1$ and $p=3$, then $\operatorname{dim}\left(\mathcal{U}_{\hat{g}}\right)=1$.
- If $q=1$ and $p>3$, then $\operatorname{dim}\left(\mathcal{U}_{\hat{g}}\right)=4 p-12$.
- If $q=p-1$, then $\operatorname{dim}\left(\mathcal{U}_{\hat{g}}\right)=3 p-7$.
- If $1<q<p-1$, then $\operatorname{dim}\left(\mathcal{U}_{\hat{g}}\right) \geq \operatorname{dim}\left(H^{1}\right)-2=-5 k-5+\sum_{i=1}^{k} 4 e_{i}$.

Consequently, if $p>2$, these spaces admit non-toric anti-self-dual deformations.
Remark 1.10. Theorem 1.8 could be equivalently stated in terms of the ALE metrics rather than the compactified metrics. However, the definition of the index on an ALE space involves defining certain weighted spaces; see [Via13, Proposition 3.1] for the precise formula which relates the index on the ALE space to the index on the compactification; for our purposes here, we only require the statement on the compactification. Similarly, Theorem 1.9 could be equivalently stated in terms of anti-self-dual ALE deformations of the ALE model.

Since the moduli space is locally given by $H^{1} / H^{0}$, it is necessary to know something about the action of $H^{0}$. This action is well known in the toric multi-Eguchi-Hanson case $q=p-1$; in this case for $p \geq 3, \operatorname{dim}\left(H^{1}\right)=3 p-6$, and the dimension of the moduli space is equal to $\operatorname{dim}\left(H^{1}\right)-1=3 p-7$. In the LeBrun negative mass case $q=1$, this action was recently completely determined by Nobuhiro Honda [Hon13] using arguments from twistor theory. For $1<q<p-1$, further arguments are needed to determine this action explicitly; this is an interesting problem.

### 1.2. Weighted projective spaces

We first recall the definition of weighted projective spaces in real dimension four:
Definition 1.11. For relatively prime integers $1 \leq r \leq q \leq p$, the weighted projective space $\mathbb{C P}_{(r, q, p)}^{2}$ is $S^{5} / S^{1}$, where $S^{1}$ acts by

$$
\begin{equation*}
\left(z_{0}, z_{1}, z_{2}\right) \mapsto\left(e^{i r \theta} z_{0}, e^{i q \theta} z_{1}, e^{i p \theta} z_{2}\right) \tag{1.14}
\end{equation*}
$$

for $0 \leq \theta<2 \pi$.
The space $\mathbb{C P}_{(r, q, p)}^{2}$ has the structure of a compact complex orbifold. Bryant [Bry01] proved that every weighted projective space admits a Bochner-Kähler metric. Subsequently, David and Gauduchon [DG06] gave a simple and direct construction of these metrics. Using an argument due to Apostolov, they also showed that this metric is the unique Bochner-Kähler metric on a given weighted projective space [DG06, Appendix D], and thus we will call this metric the canonical Bochner-Kähler metric. In complex dimension two, the Bochner tensor is the same as the anti-self-dual part of the Weyl tensor, so Bochner-Kähler metrics are the same as self-dual Kähler metrics.

The work of Derdziński [Der83] showed that a self-dual Kähler metric $g$ is conformal to a self-dual Hermitian Einstein metric on $M^{*}:=\{p \in M: R(p) \neq 0\}$, given by $\tilde{g}=R^{-2} g$, where $R$ is the scalar curvature. This conformal metric is not Kähler unless $R$ is constant. Conversely, Apostolov and Gauduchon [AG02] showed that any self-dual Hermitian Einstein metric that is not conformally flat is of the form $\tilde{g}$ for a unique selfdual Kähler metric $g$ with $R \neq 0$.

For a weighted projective space $\mathbb{C P}_{(r, q, p)}^{2}$, there are the following three cases:

- When $p<r+q$ the canonical Bochner-Kähler metric has $R>0$ everywhere, so it is conformal to a Hermitian Einstein metric with positive Einstein constant.
- When $p=r+q$ the canonical Bochner-Kähler metric has $R>0$ except at one point, so it is conformal to a complete Hermitian Einstein metric with vanishing Einstein constant outside this point.
- When $p>r+q$ the canonical Bochner-Kähler metric has $R$ vanishing along a hypersurface and the complement is composed of two open sets on which the metric is conformal to a Hermitian Einstein metric with negative Einstein constant.

For $x \in \mathbb{R},\lfloor x\rfloor$ denotes the integer part of $x$, and $\{x\}=x-\lfloor x\rfloor$ denotes the fractional part of $x$. For relatively prime integers $q$ and $r, q^{-1 ; r}$ denotes the inverse of $q$ modulo $r$. We also define the integer $\epsilon$ by

$$
\epsilon= \begin{cases}0 & \text { if } p \not \equiv q \bmod r \text { and } p \not \equiv r \bmod q,  \tag{1.15}\\ 1 & \text { if } p \equiv q \bmod r \text { or } p \equiv r \bmod q, \text { but not both, } \\ 2 & \text { if } p \equiv q \bmod r \text { and } p \equiv r \bmod q .\end{cases}
$$

Our main result for the index on weighted projective spaces is the following, with the answer depending upon certain number-theoretic properties of the triple $(r, q, p)$ :

Theorem 1.12. Let $g$ be the canonical Bochner-Kähler metric with reversed orientation on $\overline{\mathbb{C P}}_{(r, q, p)}^{2}$. Assume that $1<r<q<p$. If $r+q \geq p$ then

$$
\begin{equation*}
\operatorname{Ind}\left(\overline{\mathbb{C P}}_{(r, q, p)}^{2}, g\right)=2 \tag{1.16}
\end{equation*}
$$

If $r+q<p$, then

$$
\operatorname{Ind}\left(\overline{\mathbb{C P}}_{(r, q, p)}^{2}, g\right)= \begin{cases}2+2 \epsilon-4\left\lfloor\frac{p}{q r}\right\rfloor & \text { when }\left\{\frac{p}{q r}\right\}<\left\{\frac{q^{-1 ; r} p}{r}\right\},  \tag{1.17}\\ -2+2 \epsilon-4\left\lfloor\frac{p}{q r}\right\rfloor & \text { when }\left\{\frac{p}{q r}\right\}>\left\{\frac{q^{-1 ; r} p}{r}\right\} .\end{cases}
$$

We note that in the case $\left\{\frac{p}{q r}\right\}<\left\{\frac{q^{-1 ; r} p}{r}\right\}$, the integer $\epsilon$ can only be 0 or 1 ; the integer 2 does not actually occur in this case. Thus there are exactly five cases, which do in fact all occur (see Section 7).

Theorem 1.12 implies the following result regarding the moduli space of anti-self-dual metrics on $\overline{\mathbb{C P}}_{(r, q, p)}^{2}$ :

Theorem 1.13. Let $g$ be the canonical Bochner-Kähler metric with reversed orientation on $\overline{\mathbb{C P}}_{(r, q, p)}^{2}$. Assume that $1<r<q<p$. Then

- If $p \leq q+r$ then $[g]$ is isolated as an anti-self-dual conformal class.
- If $p>q+r$, then there exists a neighborhood, $\mathcal{U}_{g}$, of $g$ in the moduli space of anti-self-dual conformal structures with

$$
\operatorname{dim}\left(\mathcal{U}_{g}\right) \geq \begin{cases}4\left\lfloor\frac{p}{q r}\right\rfloor-2-2 \epsilon & \text { when }\left\{\frac{p}{q}\right\}<\left\{\frac{q^{-1 ; r} p}{r}\right\},  \tag{1.18}\\ 4\left\lfloor\frac{p}{q r}\right\rfloor+2-2 \epsilon & \text { when }\left\{\frac{p}{q r}\right\}>\left\{\frac{q^{-1 ; r} p}{r}\right\} .\end{cases}
$$

Remark 1.14. Since the case $p<q+r$ is conformal to an Einstein metric, it is perhaps not surprising (although not obvious) that these metrics are also isolated as self-dual metrics. But the non-trivial anti-self-dual deformations we have found in the case $p>q+r$ are quite surprising, since these metrics are rigid as Bochner-Kähler metrics.

The proof of Theorem 1.13 also relies on the fact that $H^{2}(M, g)=0$ for these metrics (see Corollary 7.10 below). Then as pointed out above, the actual moduli space is locally isomorphic to $H^{1} / H^{0}$, so the moduli space could be of dimension $\operatorname{dim}\left(H^{1}\right), \operatorname{dim}\left(H^{1}\right)-1$, or $\operatorname{dim}\left(H^{1}\right)-2$. As in the case of the Calderbank-Singer spaces, we do not determine this action explicitly here; this is another very interesting problem.

### 1.3. Outline of paper

We begin in Section 2 by recalling Kawasaki's orbifold index theorem, and apply it to the complex (1.1). Then in Section 3, we analyze the correction terms for cyclic group actions, culminating in the following formula for the index in terms of trigonometric sums when $1<q<p-1$ :

$$
\begin{align*}
\operatorname{Ind}_{\Gamma}(\hat{M})= & \frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-6+\frac{14}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \\
& -\frac{2}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \cos \left(\frac{2 \pi}{p} j\right) \cos \left(\frac{2 \pi}{p} q j\right) . \tag{1.19}
\end{align*}
$$

We note that the quantity

$$
\begin{equation*}
s(q, p)=\frac{1}{4 p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \tag{1.20}
\end{equation*}
$$

is the well-known Dedekind sum [RG72]. This has a closed form expression in several special cases, but not in general. It is not surprising that this term appears, since Dedekind sums arise naturally in the index theorem for the signature complex [HZ74, Kat87, Zag72]. However, for the anti-self-dual deformation complex, the interaction of the Dedekind sum term with the final term in (1.19) makes a huge difference. In particular, formula (1.19) says that the sum of these terms must always be an integer!

For $x \in \mathbb{R} \backslash \mathbb{Z}$, we define the sawtooth function $((x))=\{x\}-1 / 2$. In Section 4, we show that when $1<q<p-1$, the non-topological terms in (1.19) can be rewritten as a Dedekind sum plus terms involving the sawtooth function:

$$
\begin{equation*}
N(q, p)=-6+\frac{12}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)-4\left(\left(\frac{q^{-1 ; p}}{p}\right)\right)-4\left(\left(\frac{q}{p}\right)\right), \tag{1.21}
\end{equation*}
$$

where we recall that $q^{-1 ; p}$ denotes the inverse of $q$ modulo $p$. In Section 5 we use this, together with classical reciprocity for Dedekind sums, to prove Theorem 1.4. The results dealing with the Calderbank-Singer spaces, Theorems 1.8 and 1.9, are proved in Section 6. Finally, in Section 7, we present a complete analysis of the index for the canonical Bochner-Kähler metric on a weighted projective space, and prove Theorem 1.13. Interestingly, an important ingredient is Rademacher's triple reciprocity formula for Dedekind sums [Rad54]. We conclude the paper with some remarks on the number-theoretic condition on the triple $(r, q, p)$ which occurs in Theorem 1.12.

## 2. The orbifold $\Gamma$-index

For an orbifold $(M, g)$, the $\Gamma$-index is given analytically by

$$
\begin{equation*}
\operatorname{Ind}_{\Gamma}(M, g)=\operatorname{dim}\left(H^{0}(M, g)\right)-\operatorname{dim}\left(H^{1}(M, g)\right)+\operatorname{dim}\left(H^{2}(M, g)\right) \tag{2.1}
\end{equation*}
$$

From Kawasaki's orbifold index theorem [Kaw81], it follows that we have a $\Gamma$-index formula of the form

$$
\begin{equation*}
\operatorname{Ind}_{\Gamma}(M)=\frac{1}{2}\left(15 \chi_{\text {orb }}(M)+29 \tau_{\text {orb }}(M)\right)+\frac{1}{|\Gamma|} \sum_{\gamma \neq \mathrm{ld}} \frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)} \tag{2.2}
\end{equation*}
$$

where $\chi_{\text {orb }}(M)$ is the orbifold Euler characteristic defined by

$$
\begin{equation*}
\chi_{\mathrm{orb}}(M)=\frac{1}{8 \pi^{2}} \int_{M}\left(|W|^{2}-\frac{1}{2}|\operatorname{Ric}|^{2}+\frac{1}{6} R^{2}\right) d V_{g}, \tag{2.3}
\end{equation*}
$$

the quantity $\tau_{\text {orb }}(M)$ is the orbifold signature defined by

$$
\begin{equation*}
\tau_{\mathrm{orb}}(M)=\frac{1}{12 \pi^{2}} \int_{M}\left(\left|W^{+}\right|^{2}-\left|W^{-}\right|^{2}\right) d V_{g} \tag{2.4}
\end{equation*}
$$

and the quantity $\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right) / \operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)$ is a correction term depending upon the action of $\gamma$ on certain bundles, which we will describe in what follows.

In the next subsection, we compute the trace of the action of $\gamma$, an element in the orbifold group $\Gamma$, on the bundles $\left[N_{\mathbb{C}}\right],\left[S_{0}^{2}\left(N_{\mathbb{C}}\right)\right]$ and $\left[S_{0}^{2}\left(\Lambda_{+}^{2}\right)\right]$ over the fixed point set, which we then use to compute a general formula for the $\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right) / \mathrm{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)$ term. Then we give the orbifold Euler characteristic and orbifold signature in terms of the topological Euler characteristic and topological signature and correction terms also depending upon the $\gamma$-action respectively. Finally, we combine this information into a formula for the orbifold $\Gamma$-index.

### 2.1. Group action on bundles

In order to compute the $\Gamma$-index, we first need to find the trace of the $\gamma$-action, for every $\gamma$ in $\Gamma$, on the pullback of the complexified principal symbol, $i^{*} \sigma$, where

$$
\begin{equation*}
i: p \rightarrow M \tag{2.5}
\end{equation*}
$$

is the inclusion map from the fixed point $p$ into the orbifold $M$. In this case

$$
\begin{equation*}
i^{*} \sigma=\left[N_{\mathbb{C}}\right]-\left[S_{0}^{2}\left(N_{\mathbb{C}}\right)\right]+\left[S_{0}^{2} \Lambda_{+}^{2}\right] \tag{2.6}
\end{equation*}
$$

For a general $\gamma$ of the form

$$
\gamma=\left(\begin{array}{cccc}
\cos \theta_{1} & -\sin \theta_{1} & 0 & 0  \tag{2.7}\\
\sin \theta_{1} & \cos \theta_{1} & 0 & 0 \\
0 & 0 & \cos \theta_{2} & -\sin \theta_{2} \\
0 & 0 & \sin \theta_{2} & \cos \theta_{2}
\end{array}\right)
$$

fixing the point $p$, the normal bundle is trivial, so $N_{\mathbb{C}}:=N \otimes \mathbb{C}=\mathbb{C}^{4}$, and we have the following proposition.

Proposition 2.1. The trace of the $\gamma$-action on the components of $i^{*} \sigma$ is as follows:
(1) $\operatorname{tr}\left(\left.\gamma\right|_{N_{\mathbb{C}}}\right)=2 \cos \left(\theta_{1}\right)+2 \cos \left(\theta_{2}\right)$,
(2) $\operatorname{tr}\left(\left.\gamma\right|_{S_{0}^{2}\left(N_{\mathbb{C}}\right)}\right)=1+2 \cos \left(\theta_{1}+\theta_{2}\right)+2 \cos \left(-\theta_{1}+\theta_{2}\right)+4 \cos \left(\theta_{1}+\theta_{2}\right) \cos \left(-\theta_{1}+\theta_{2}\right)$,
(3) $\operatorname{tr}\left(\left.\gamma\right|_{S_{0}^{2}\left(\Lambda_{+}^{2}\right)}\right)=2 \cos \left(\theta_{1}+\theta_{2}\right)+4 \cos ^{2}\left(\theta_{1}+\theta_{2}\right)-1$.

Proof. The normal bundle can be written as $N=x_{1} \oplus \cdots \oplus x_{4}$ in real coordinates. After complexifying the normal bundle we can diagonalize $\gamma$ to write

$$
\left.\gamma\right|_{N_{\mathbb{C}}}=\left(\begin{array}{cccc}
e^{i \theta_{1}} & 0 & 0 & 0  \tag{2.8}\\
0 & e^{-i \theta_{1}} & 0 & 0 \\
0 & 0 & e^{i \theta_{2}} & 0 \\
0 & 0 & 0 & e^{-i \theta_{2}}
\end{array}\right)
$$

with respect to the complex basis $\left\{\lambda_{1} \oplus \lambda_{2} \oplus \lambda_{3} \oplus \lambda_{4}\right\}=\mathbb{C}^{4}$, where

$$
\begin{equation*}
\left\{2 x_{1}, 2 x_{2}, 2 x_{3}, 2 x_{4}\right\}=\left\{\lambda_{1}-i \lambda_{2}, i \lambda_{1}-\lambda_{2}, \lambda_{3}-i \lambda_{4}, i \lambda_{3}-\lambda_{4}\right\} . \tag{2.9}
\end{equation*}
$$

Formula (1) follows immediately.
Next, to see how $\gamma$ acts on $S_{0}^{2}\left(N_{g}\right)=\Lambda_{+}^{2} \otimes \Lambda_{-}^{2}$ we first examine how $\gamma$ acts on $\Lambda_{+}^{2}$ and $\Lambda_{-}^{2}$ independently. We use the following basis for $\Lambda_{+}^{2}$ :

$$
\begin{align*}
& \omega_{1}^{+}=\frac{1}{2}\left[d \lambda_{2} \wedge d \lambda_{1}+d \lambda_{4} \wedge d \lambda_{3}\right] \\
& \omega_{2}^{+}=\frac{1}{2}\left[d \lambda_{1} \wedge d \lambda_{3}+d \lambda_{4} \wedge d \lambda_{2}\right]  \tag{2.10}\\
& \omega_{3}^{+}=\frac{1}{2}\left[i d \lambda_{1} \wedge d \lambda_{3}+i d \lambda_{2} \wedge d \lambda_{4}\right]
\end{align*}
$$

and the following basis for $\Lambda_{-}^{2}$ :

$$
\begin{align*}
& \omega_{1}^{-}=\frac{1}{2}\left[d \lambda_{2} \wedge d \lambda_{1}-d \lambda_{4} \wedge d \lambda_{3}\right] \\
& \omega_{2}^{-}=\frac{1}{2}\left[i d \lambda_{3} \wedge d \lambda_{2}+i d \lambda_{4} \wedge d \lambda_{1}\right]  \tag{2.11}\\
& \omega_{3}^{-}=\frac{1}{2}\left[d \lambda_{2} \wedge d \lambda_{3}+d \lambda_{4} \wedge d \lambda_{1}\right]
\end{align*}
$$

So we see that $\gamma$ acts on $\Lambda_{+}^{2}$ by

$$
\begin{align*}
& \gamma\left(\omega_{1}^{+}\right)=\omega_{1}^{+}, \\
& \gamma\left(\omega_{2}^{+}\right)=\frac{1}{2}\left[e^{i\left(\theta_{1}+\theta_{2}\right)}\left(\omega_{2}^{+}-i \omega_{3}^{+}\right)+e^{-i\left(\theta_{1}+\theta_{2}\right)}\left(\omega_{2}^{+}+i \omega_{3}^{+}\right)\right],  \tag{2.12}\\
& \gamma\left(\omega_{3}^{+}\right)=\frac{1}{2}\left[e^{i\left(\theta_{1}+\theta_{2}\right)}\left(\omega_{3}^{+}+i \omega_{2}^{+}\right)+e^{-i\left(\theta_{1}+\theta_{2}\right)}\left(\omega_{3}^{+}-i \omega_{2}^{+}\right)\right],
\end{align*}
$$

and $\gamma$ acts on $\Lambda_{-}^{2}$ by

$$
\begin{align*}
& \gamma\left(\omega_{1}^{-}\right)=\omega_{1}^{-}, \\
& \gamma\left(\omega_{2}^{-}\right)=\frac{1}{2}\left[e^{i\left(-\theta_{1}+\theta_{2}\right)}\left(\omega_{2}^{-}-i \omega_{3}^{-}\right)+e^{i\left(\theta_{1}-\theta_{2}\right)}\left(\omega_{2}^{-}+i \omega_{3}^{-}\right)\right],  \tag{2.13}\\
& \gamma\left(\omega_{3}^{+}\right)=\frac{1}{2}\left[e^{i\left(-\theta_{1}+\theta_{2}\right)}\left(\omega_{3}^{-}+i \omega_{2}^{-}\right)+e^{i\left(\theta_{1}-\theta_{2}\right)}\left(\omega_{3}^{-}-i \omega_{2}^{-}\right)\right] .
\end{align*}
$$

Therefore, we can write

$$
\begin{align*}
&\left.\gamma\right|_{\Lambda_{+}^{2}}=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & \cos \left(\theta_{1}+\theta_{2}\right) & -\sin \left(\theta_{1}+\theta_{2}\right) \\
0 & \sin \left(\theta_{1}+\theta_{2}\right) & \cos \left(\theta_{1}+\theta_{2}\right)
\end{array}\right)  \tag{2.14}\\
&\left.\gamma\right|_{\Lambda_{-}^{2}}=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & \cos \left(-\theta_{1}+\theta_{2}\right) & -\sin \left(-\theta_{1}+\theta_{2}\right) \\
0 & \sin \left(-\theta_{1}+\theta_{2}\right) & \cos \left(-\theta_{1}+\theta_{2}\right)
\end{array}\right) . \tag{2.15}
\end{align*}
$$

To derive (2), we compute

$$
\begin{align*}
& \operatorname{tr}\left(\left.\gamma\right|_{S_{0}^{2} N_{\mathbb{C}}}\right)=\operatorname{tr}\left(\left.\gamma\right|_{\Lambda_{+}^{2} \otimes \Lambda_{-}^{2}}\right)=\operatorname{tr}\left(\left.\gamma\right|_{\Lambda_{+}^{2}}\right) \cdot \operatorname{tr}\left(\left.\gamma\right|_{\Lambda_{-}^{2}}\right) \\
& \quad=\left(1+2 \cos \left(\theta_{1}+\theta_{2}\right)\right) \cdot\left(1+2 \cos \left(-\theta_{1}+\theta_{2}\right)\right) \\
& \quad=1+2 \cos \left(\theta_{1}+\theta_{2}\right)+2 \cos \left(-\theta_{1}+\theta_{2}\right)+4 \cos \left(\theta_{1}+\theta_{2}\right) \cos \left(-\theta_{1}+\theta_{2}\right) . \tag{2.16}
\end{align*}
$$

Next, to see how $\gamma$ acts on $S_{0}^{2}\left(\Lambda_{+}^{2}\right)$, decompose

$$
\begin{equation*}
S_{0}^{2} \Lambda_{+}^{2}=\left[\mathbb{C} \otimes\left(\omega_{2}^{+} \oplus \omega_{3}^{+}\right)\right] \oplus S_{0}^{2}\left(\omega_{2}^{+} \oplus \omega_{3}^{+}\right) \oplus t r \tag{2.17}
\end{equation*}
$$

where $\operatorname{tr}=2 \omega_{1}^{+}-\left(\omega_{2}^{+}+\omega_{3}^{+}\right)$denotes the trace component, and write the basis of $S_{0}^{2}\left(\omega_{2}^{+} \oplus \omega_{3}^{+}\right)$as

$$
\begin{equation*}
\left\{\omega_{2}^{+} \otimes \omega_{2}^{+}-\omega_{3}^{+} \otimes \omega_{3}^{+}, \omega_{2}^{+} \otimes \omega_{3}^{+}+\omega_{3}^{+} \otimes \omega_{2}^{+}\right\} \tag{2.18}
\end{equation*}
$$

We see that

$$
\left.\gamma\right|_{\omega_{1}^{+} \otimes\left(\omega_{2}^{+} \oplus \omega_{3}^{+}\right)}=\left(\begin{array}{cc}
\cos \left(\theta_{1}+\theta_{2}\right) & -\sin \left(\theta_{1}+\theta_{2}\right)  \tag{2.19}\\
\sin \left(\theta_{1}+\theta_{2}\right) & \cos \left(\theta_{1}+\theta_{2}\right)
\end{array}\right)
$$

$$
\begin{align*}
\left.\gamma\right|_{S_{0}^{2}\left(\omega_{2}^{+} \oplus \omega_{3}^{2}\right)} & =\left(\begin{array}{cc}
\cos ^{2}\left(\theta_{1}+\theta_{2}\right)-\sin ^{2}\left(\theta_{1}+\theta_{2}\right) & -2 \sin \left(\theta_{1}+\theta_{2}\right) \cos \left(\theta_{1}+\theta_{2}\right) \\
2 \sin \left(\theta_{1}+\theta_{2}\right) \cos \left(\theta_{1}+\theta_{2}\right) & \cos ^{2}\left(\theta_{1}+\theta_{2}\right)-\sin ^{2}\left(\theta_{1}+\theta_{2}\right)
\end{array}\right),  \tag{2.20}\\
\left.\gamma\right|_{t r \in S_{0}^{2} \Lambda_{+}^{2}} & =1 . \tag{2.21}
\end{align*}
$$

Using these, we derive (3) by computing

$$
\begin{align*}
\operatorname{tr}\left(\left.\gamma\right|_{S_{0}^{2} \Lambda_{+}^{2}}\right) & =\left[2 \cos \left(\theta_{1}+\theta_{2}\right)\right]+\left[4 \cos ^{2}\left(\theta_{1}+\theta_{2}\right)-2\right]+[1] \\
& =2 \cos \left(\theta_{1}+\theta_{2}\right)+4 \cos ^{2}\left(\theta_{1}+\theta_{2}\right)-1 \tag{2.22}
\end{align*}
$$

### 2.2. Equivariant Chern character

We next compute the term $\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right) / \operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)$. The numerator of this term is the $\gamma$-equivariant Chern character of the pullback of the principal symbol, $i^{*} \sigma$, described in (2.5) and (2.6). The denominator is the $\gamma$-equivariant Chern character of the $K$-theoretic Thom class of the complexified normal bundle. Since the normal bundle is trivial over the fixed point, this is

$$
\begin{equation*}
\lambda_{-1} N_{\mathbb{C}}=\left[\Lambda^{0}\left(\mathbb{C}^{4}\right)\right]-\left[\Lambda^{1}\left(\mathbb{C}^{4}\right)\right]+\left[\Lambda^{2}\left(\mathbb{C}^{4}\right)\right]-\left[\Lambda^{3}\left(\mathbb{C}^{4}\right)\right]+\left[\Lambda^{4}\left(\mathbb{C}^{4}\right)\right] \tag{2.23}
\end{equation*}
$$

Since the $\gamma$-equivariant Chern character is just the $\gamma$-action times the Chern character of each eigenspace, using Proposition 2.1, we compute

$$
\begin{align*}
\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)= & \operatorname{tr}\left(\left.\gamma\right|_{N_{\mathbb{C}}}\right)-\operatorname{tr}\left(\left.\gamma\right|_{S_{0}^{2} N_{\mathbb{C}}}\right)+\operatorname{tr}\left(\left.\gamma\right|_{S_{0}^{2} \Lambda_{+}^{2}}\right) \\
= & {\left[2 \cos \theta_{1}+2 \cos \theta_{2}\right] } \\
& -\left[1+2 \cos \left(\theta_{1}+\theta_{2}\right)+2 \cos \left(-\theta_{1}+\theta_{2}\right)+4 \cos \left(\theta_{1}+\theta_{2}\right) \cos \left(-\theta_{1}+\theta_{2}\right)\right] \\
& +\left[2 \cos \left(\theta_{1}+\theta_{2}\right)+4 \cos ^{2}\left(\theta_{1}+\theta_{2}\right)-1\right] \\
= & {\left[2 \cos \theta_{1}+2 \cos \theta_{2}-2-2 \cos \left(\theta_{1}\right) \cos \left(\theta_{2}\right)\right] } \\
& +\left[-2 \sin \left(\theta_{1}\right) \sin \left(\theta_{2}\right)-8 \cos \left(\theta_{1}\right) \cos \left(\theta_{2}\right) \sin \left(\theta_{1}\right) \sin \left(\theta_{2}\right)+8 \sin ^{2}\left(\theta_{1}\right) \sin ^{2}\left(\theta_{2}\right)\right] \\
= & {\left[-2\left(\cos \theta_{1}-1\right)\left(\cos \theta_{2}-1\right)\right]+\left[8\left(1-\cos ^{2} \theta_{1}\right)\left(1-\cos ^{2} \theta_{2}\right)\right] } \\
& +\left[-2 \sin \left(\theta_{1}\right) \sin \left(\theta_{2}\right)-8 \cos \left(\theta_{1}\right) \cos \left(\theta_{2}\right) \sin \left(\theta_{1}\right) \sin \left(\theta_{2}\right)\right] . \tag{2.24}
\end{align*}
$$

Similarly, we compute

$$
\begin{align*}
\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)= & \operatorname{tr}\left(\left.\gamma\right|_{\left[\Lambda^{0}\left(\mathbb{C}^{4}\right)\right]}\right)-\operatorname{tr}\left(\left.\gamma\right|_{\left[\Lambda^{1}\left(\mathbb{C}^{4}\right)\right]}\right)+\operatorname{tr}\left(\left.\gamma\right|_{\left[\Lambda^{2}\left(\mathbb{C}^{4}\right)\right]}\right) \\
& -\operatorname{tr}\left(\left.\gamma\right|_{\left[\Lambda^{3}\left(\mathbb{C}^{4}\right)\right]}\right)+\operatorname{tr}\left(\left.\gamma\right|_{\left[\Lambda^{4}\left(\mathbb{C}^{4}\right)\right]}\right)=4\left(\cos \theta_{1}-1\right)\left(\cos \theta_{2}-1\right) . \tag{2.25}
\end{align*}
$$

Therefore

$$
\begin{align*}
\frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)}= & {\left[-\frac{1}{2}+2\left(1+\cos \theta_{1}\right)\left(1+\cos \theta_{2}\right)\right] } \\
& -\left[\frac{2 \sin \left(\theta_{1}\right) \sin \left(\theta_{2}\right)+8 \cos \left(\theta_{1}\right) \cos \left(\theta_{2}\right) \sin \left(\theta_{1}\right) \sin \left(\theta_{2}\right)}{4\left(\cos \theta_{1}-1\right)\left(\cos \theta_{2}-1\right)}\right] \tag{2.26}
\end{align*}
$$

Since $\frac{\sin \left(\theta_{1}\right) \sin \left(\theta_{2}\right)}{\left(\cos \theta_{1}-1\right)\left(\cos \theta_{2}-1\right)}=\cot \left(\frac{\theta_{1}}{2}\right) \cot \left(\frac{\theta_{2}}{2}\right)$, we see that

$$
\begin{align*}
\frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)}= & -\frac{1}{2}+2\left(1+\cos \theta_{1}\right)\left(1+\cos \theta_{2}\right)-\frac{1}{2} \cot \left(\frac{\theta_{1}}{2}\right) \cot \left(\frac{\theta_{2}}{2}\right) \\
& -2 \cot \left(\frac{\theta_{1}}{2}\right) \cot \left(\frac{\theta_{2}}{2}\right) \cos \left(\theta_{1}\right) \cos \left(\theta_{2}\right) \tag{2.27}
\end{align*}
$$

### 2.3. The $\Gamma$-index

For an orbifold with a single isolated singularity, we have a formula for the Euler characteristic,

$$
\begin{equation*}
\chi_{\mathrm{top}}(M)=\chi_{\mathrm{orb}}(M)+\frac{|\Gamma|-1}{|\Gamma|} \tag{2.28}
\end{equation*}
$$

and a formula for the signature,

$$
\begin{equation*}
\tau_{\text {top }}(M)=\tau_{\text {orb }}(M)-\eta\left(S^{3} / \Gamma\right) \tag{2.29}
\end{equation*}
$$

where $\Gamma \subset \mathrm{SO}(4)$ is the orbifold group around the fixed point and $\eta\left(S^{3} / \Gamma\right)$ is the etainvariant, which in our case is given by

$$
\begin{equation*}
\eta\left(S^{3} / \Gamma\right)=\frac{1}{|\Gamma|} \sum_{\gamma \neq \mathrm{ld}} \cot \left(\frac{\theta_{1}}{2} j\right) \cot \left(\frac{\theta_{2}}{2} j\right) \tag{2.30}
\end{equation*}
$$

See [Hit97] for a useful discussion of the formulas (2.28) and (2.29).
Combining (2.28) and (2.29) with the formula for the $\Gamma$-index given in (2.2), we have

$$
\begin{equation*}
\operatorname{Ind}_{\Gamma}=\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-\frac{15}{2}\left(\frac{|\Gamma|-1}{|\Gamma|}\right)+\frac{29}{2} \eta\left(S^{3} / \Gamma\right)+\frac{1}{|\Gamma|} \sum_{\gamma \neq \mathrm{ld}} \frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)} \tag{2.31}
\end{equation*}
$$

where the last term is given by formula (2.27).

## 3. $\Gamma$-index for cyclic group actions

We consider an orbifold with an isolated singularity having the group action $\Gamma_{(q, p)}$ generated by

$$
\gamma=\left(\begin{array}{cccc}
\cos \left(\frac{2 \pi}{p}\right) & -\sin \left(\frac{2 \pi}{p}\right) & 0 & 0  \tag{3.1}\\
\sin \left(\frac{2 \pi}{p}\right) & \cos \left(\frac{2 \pi}{p}\right) & 0 & 0 \\
0 & 0 & \cos \left(\frac{2 \pi}{p} q\right) & -\sin \left(\frac{2 \pi}{p} q\right) \\
0 & 0 & \sin \left(\frac{2 \pi}{p} q\right) & \cos \left(\frac{2 \pi}{p} q\right)
\end{array}\right),
$$

where $p$ and $q$ are relatively prime. The cases when $q=1$ and $q=p-1$ have already been resolved in [Via13], and although we are specifically interested in $1<q<p-1$,
we will make use of the sum

$$
\begin{equation*}
\sum_{\gamma \neq \mathrm{ld}} \frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)} \tag{3.2}
\end{equation*}
$$

in all cases, and make our computations accordingly. We begin this section by simplifying our formula for this sum in general:

$$
\begin{align*}
\sum_{\gamma \neq \mathrm{ld}} & \frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)} \\
= & \sum_{j=1}^{p-1}\left[-\frac{1}{2}+2\left(1+\cos \left(\frac{2 \pi}{p} j\right)\right)\left(1+\cos \left(\frac{2 \pi}{p} q j\right)\right)-\frac{1}{2} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)\right] \\
& \quad-\sum_{j=1}^{p-1} 2 \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \cos \left(\frac{2 \pi}{p} j\right) \cos \left(\frac{2 \pi}{p} q j\right) \\
= & \sum_{j=1}^{p-1}\left[\frac{3}{2}+2 \cos \left(\frac{2 \pi}{p} j\right)+2 \cos \left(\frac{2 \pi}{p} q j\right)+\cos \left(\frac{2 \pi}{p}(q+1) j\right)\right] \\
& +\sum_{j=1}^{p-1}\left[\cos \left(\frac{2 \pi}{p}(q-1) j\right)-\frac{1}{2} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)\right] \\
& +\sum_{j=1}^{p-1}\left[-2 \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \cos \left(\frac{2 \pi}{p} j\right) \cos \left(\frac{2 \pi}{p} q j\right)\right] \tag{3.3}
\end{align*}
$$

Now, to further simplify our formula for the $\Gamma$-index, it is necessary to separate into the following cases:

## 3.1. $\Gamma$-index when $1<q<p-1$

Using (3.3), we see that in this case

$$
\begin{align*}
\sum_{\gamma \neq \mathrm{ld}} \frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)}= & {\left[\frac{3}{2} p-\frac{15}{2}\right]-\frac{1}{2} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) } \\
& -2 \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \cos \left(\frac{2 \pi}{p} j\right) \cos \left(\frac{2 \pi}{p} q j\right) \tag{3.4}
\end{align*}
$$

Therefore, by combining this with formula (2.31) for the $\Gamma$-index, we have

$$
\begin{align*}
\operatorname{Ind}_{\Gamma}(M)= & \frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-6+\frac{14}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \\
& -\frac{2}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \cos \left(\frac{2 \pi}{p} j\right) \cos \left(\frac{2 \pi}{p} q j\right) \tag{3.5}
\end{align*}
$$

3.2. $\Gamma$-index when $q=1$ and $p=2$

Using (3.3), we see that in this case

$$
\begin{equation*}
\sum_{\gamma \neq \mathrm{ld}} \frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)}=-\frac{1}{2} \tag{3.6}
\end{equation*}
$$

Therefore, by combining this with (2.31), we have

$$
\begin{equation*}
\operatorname{Ind}_{\Gamma}(M)=\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-4 \tag{3.7}
\end{equation*}
$$

3.3. $\Gamma$-index when $q=1$ and $p>2$

Using (3.3), we see that in this case

$$
\begin{equation*}
\sum_{\gamma \neq \mathrm{Id}} \frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{\gamma}\left(\lambda_{-1} N_{\mathbb{C}}\right)}=\frac{5}{2} p-\frac{15}{2}-\sum_{j=1}^{p-1}\left[\frac{1}{2} \cot ^{2}\left(\frac{\pi}{p} j\right)+2 \cot ^{2}\left(\frac{\pi}{p} j\right) \cos ^{2}\left(\frac{2 \pi}{p} j\right)\right] \tag{3.8}
\end{equation*}
$$

Therefore, by combining this with (2.31), and the following well known formula for the Dedekind sum (see [RG72]):

$$
\begin{equation*}
\frac{1}{4 p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right)=\frac{1}{12 p}(p-1)(p-2) \tag{3.9}
\end{equation*}
$$

we have

$$
\begin{align*}
\operatorname{Ind}_{\Gamma}(M) & =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-5+\frac{14}{p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right)-\frac{2}{p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right) \cos ^{2}\left(\frac{2 \pi}{p} j\right) \\
& =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-5+\frac{12}{p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right)+\frac{8}{p} \sum_{j=1}^{p-1} \cos ^{4}\left(\frac{\pi}{p} j\right) \\
& =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-2-\frac{8}{p}+\frac{12}{p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right) \\
& =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-2-\frac{8}{p}+\frac{4}{p}\left(p^{2}-3 p+2\right) \\
& =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)+4 p-14 \tag{3.10}
\end{align*}
$$

## 3.4. $\Gamma$-index when $q=p-1$ and $p>2$

Using (3.3), we see that in this case

$$
\begin{equation*}
\sum_{\gamma \neq \mathrm{Id}} \frac{\operatorname{ch}_{\gamma}\left(i^{*} \sigma\right)}{\operatorname{ch}_{g}\left(\lambda_{-1} N_{\mathbb{C}}\right)}=\frac{5}{2} p-\frac{15}{2}+\sum_{j=1}^{p-1}\left[\frac{1}{2} \cot ^{2}\left(\frac{\pi}{p} j\right)+2 \cot ^{2}\left(\frac{\pi}{p} j\right) \cos ^{2}\left(\frac{2 \pi}{p} j\right)\right] \tag{3.11}
\end{equation*}
$$

Therefore, by combining this with (2.31), we have

$$
\begin{align*}
\operatorname{Ind}_{\Gamma}(M) & =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-5-\frac{14}{p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right)+\frac{2}{p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right) \cos ^{2}\left(\frac{2 \pi}{p} j\right) \\
& =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-8+\frac{8}{p}-\frac{12}{p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right) \\
& =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-8+\frac{8}{p}-\frac{4}{p}\left(p^{2}-3 p+2\right) \\
& =\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)-4 p+4 \tag{3.12}
\end{align*}
$$

## 4. Non-topological terms in the $\Gamma$-index

We denote the terms in the $\Gamma$-index not involving the topological Euler characteristic or topological signature by $N(q, p)$. Also we change our notation of the $\Gamma$-index from $\operatorname{Ind}_{\Gamma}$ to $\operatorname{Ind}_{(q, p)}$ to reflect the particular group action. With this new notation we can write the index as

$$
\begin{equation*}
\operatorname{Ind}_{(q, p)}=\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)+N(q, p) \tag{4.1}
\end{equation*}
$$

In this section we will simplify our formulas for $N(q, p)$. Also, for the remainder of the paper we will use the following notation. For two relatively prime positive integers $\alpha<\beta$, denote $\alpha$ 's inverse modulo $\beta$ by $\alpha^{-1 ; \beta}$, and $\beta$ 's inverse modulo $\alpha$ by $\beta^{-1 ; \alpha}$, i.e.

$$
\begin{equation*}
\alpha \alpha^{-1 ; \beta} \equiv 1 \bmod \beta \quad \text { and } \quad \beta \beta^{-1 ; \alpha} \equiv 1 \bmod \alpha . \tag{4.2}
\end{equation*}
$$

In the cases that $N(q, p)$ is easy to compute we see that

$$
N(q, p)= \begin{cases}4 p-14 & \text { when } 1=q<p-1  \tag{4.3}\\ -4 p+4 & \text { when } q=p-1\end{cases}
$$

Note that the case when $q= \pm 1$ and $p=2$ can be actually included in the $q=p-1$ case. It will be convenient later if we also have these formulas written in terms of sawtooth functions, a cotangent sum and a constant where the sawtooth function is defined to be

$$
((x))= \begin{cases}x-\lfloor x\rfloor-1 / 2 & \text { when } x \notin \mathbb{Z}  \tag{4.4}\\ 0 & \text { when } x \in \mathbb{Z} .\end{cases}
$$

We will include the formulas from (4.3), written in this way, below in Theorem 4.2.
To compute $N(q, p)$ in all other cases we will employ the following proposition:

## Proposition 4.1.

$$
\begin{equation*}
-\frac{1}{2 p} \sum_{j=1}^{p-1} \sin \left(\frac{2 \pi}{p} q j\right) \cot \left(\frac{\pi}{p} j\right)=\left(\left(\frac{q}{p}\right)\right) \tag{4.5}
\end{equation*}
$$

which is the sawtooth function defined in (4.4).
Proof. This is due to Eisenstein; see [Apo90].

Theorem 4.2. When $q \not \equiv(p-1) \bmod p$ we have the formula

$$
\begin{equation*}
N(q, p)=-6+\frac{12}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)-4\left(\left(\frac{q^{-1 ; p}}{p}\right)\right)-4\left(\left(\frac{q}{p}\right)\right) \tag{4.6}
\end{equation*}
$$

and when $q \equiv(p-1) \bmod p$ we have

$$
\begin{equation*}
N(q, p)=N(p-1, p)=-4-\frac{12}{p} \sum_{j=1}^{p-1} \cot ^{2}\left(\frac{\pi}{p} j\right)+8\left(\left(\frac{1}{p}\right)\right) \tag{4.7}
\end{equation*}
$$

Proof. For the $q=p-1$ case, by examining the formulas in (3.10) and (3.12), one can easily see that we can also write $N(p-1, p)=-4 p+4$ in this way. Now, consider the $1 \leq q<p$ case. From (3.5), we begin by computing

$$
\begin{align*}
N(q, p)= & -6+\frac{14}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \\
& -\frac{2}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \cos \left(\frac{2 \pi}{p} j\right) \cos \left(\frac{2 \pi}{p} q j\right) \\
= & -6+\frac{2}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)\left[7-\cos \left(\frac{2 \pi}{p} j\right) \cos \left(\frac{2 \pi}{p} q j\right)\right] \tag{4.8}
\end{align*}
$$

and using the identity $\cos (2 x)=1-2 \sin ^{2}(x)$, we expand this to

$$
\begin{align*}
&-6+\frac{2}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)\left[7-\left(1-2 \sin ^{2}\left(\frac{\pi}{p} j\right)\right)\left(1-2 \sin ^{2}\left(\frac{\pi}{p} q j\right)\right)\right] \\
&=-6+\frac{2}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)\left[6+2 \sin ^{2}\left(\frac{\pi}{p} j\right)+2 \sin ^{2}\left(\frac{\pi}{p} q j\right)\right] \\
&+\frac{2}{p} \sum_{j=1}^{p-1} \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)\left[-4 \sin ^{2}\left(\frac{\pi}{p} j\right) \sin ^{2}\left(\frac{\pi}{p} q j\right)\right] \tag{4.9}
\end{align*}
$$

which simplifies to

$$
\begin{align*}
N(q, p)= & -6+\frac{1}{p} \sum_{j=1}^{p-1} 12 \cot \left(\frac{\pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right)+\frac{1}{p} \sum_{j=1}^{p-1} 2 \sin \left(\frac{2 \pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) \\
& +\frac{1}{p} \sum_{j=1}^{p-1} 4 \sin \left(\frac{\pi}{p} q j\right) \cos \left(\frac{\pi}{p} q j\right) \cot \left(\frac{\pi}{p} j\right) \\
& -\frac{1}{p} \sum_{j=1}^{p-1} 8 \sin \left(\frac{\pi}{p} j\right) \cos \left(\frac{\pi}{p} j\right) \sin \left(\frac{\pi}{p} q j\right) \cos \left(\frac{\pi}{p} q j\right) \tag{4.10}
\end{align*}
$$

The fifth term on the right hand side of (4.10) sums to zero because

$$
\begin{gather*}
\frac{-8}{p} \sum_{j=1}^{p-1} \sin \left(\frac{\pi}{p} j\right) \cos \left(\frac{\pi}{p} j\right) \sin \left(\frac{\pi}{p} q j\right) \cos \left(\frac{\pi}{p} q j\right)=\frac{-4}{p} \sum_{j=1}^{p-1} \sin \left(\frac{2 \pi}{p} j\right) \sin \left(\frac{2 \pi}{p} q j\right) \\
=\frac{-2}{p} \sum_{j=1}^{p-1}\left[\cos \left(\frac{2 \pi}{p}(1-q) j\right)-\cos \left(\frac{2 \pi}{p}(1+q) j\right)\right]=0 . \tag{4.11}
\end{gather*}
$$

By Proposition 4.1, the fourth term on the right hand side of (4.10) is

$$
\frac{4}{p} \sum_{j=1}^{p-1} \sin \left(\frac{\pi}{p} q j\right) \cos \left(\frac{\pi}{p} q j\right) \cot \left(\frac{\pi}{p} j\right)=\frac{2}{p} \sum_{j=1}^{p-1} \sin \left(2 \frac{\pi}{p} q j\right) \cot \left(\frac{\pi}{p} j\right)=-4\left(\left(\frac{q}{p}\right)\right),
$$

and the third term on the right hand side of (4.10) is

$$
\begin{aligned}
\frac{2}{p} \sum_{j=1}^{p-1} \sin \left(\frac{2 \pi}{p} j\right) \cot \left(\frac{\pi}{p} q j\right) & =\frac{2}{p} \sum_{j=1}^{p-1} \sin \left(\frac{2 \pi}{p} q q^{-1 ; p} j\right) \cot \left(\frac{\pi}{p} q j\right) \\
& =\frac{2}{p} \sum_{r=1}^{p-1} \sin \left(\frac{2 \pi}{p} q^{-1 ; p} r\right) \cot \left(\frac{\pi}{p} r\right)=-4\left(\left(\frac{q^{-1 ; p}}{p}\right)\right)
\end{aligned}
$$

where $r=j q^{-1 ; p}$; this finishes the proof.
Since the formulas for $N(q, p)$ given in Theorem 4.2 are the same in all cases except when $q=p-1$, we make the following definition:

Definition 4.3. A singularity is said to be exceptional if it results from a $(p-1, p)$-action.
Otherwise, it is called non-exceptional.

## 5. Explicit formula for $N(q, p)$

We begin this section by proving reciprocity formulas for the individual summands of $N(q, p)$. Then, we use these relations to prove reciprocity formulas for $N(q, p)$, which will later be used to compute $N(q, p)$ explicitly. Since we have already computed $N(1, p)$, for the simplicity of presentation, we will assume that $q>1$ for the following. To simplify notation we let $A(q, p)=48 s(q, p)$, where $s(q, p)$ is the Dedekind sum defined in (1.20).
Proposition 5.1. Writing $p=e q-a$, we have the following reciprocity relations:
(1) $A(q, p)+A(p, q)=-12+4 e-4 \frac{a}{q}+4 \frac{q}{p}+4 \frac{1}{p q}$,
(2) $-4\left(\left(\frac{q^{-1 ; p}}{p}\right)\right)-4\left(\left(\frac{p^{-1 ; q}}{q}\right)\right)=-\frac{4}{p q}$,
(3) $-4\left(\left(\frac{q}{p}\right)\right)-4\left(\left(\frac{p}{q}\right)\right)=-4 \frac{q}{p}+4 \frac{a}{q}$.

Proof. By the reciprocity formula for Dedekind sums [RG72], we have

$$
\left.\begin{array}{rl}
A(q, p)+A & (p, q)
\end{array}\right)=-12+4\left(\frac{p}{q}+\frac{q}{p}+\frac{1}{p q}\right) .
$$

Next,

$$
\begin{align*}
-4\left(\left(\frac{q^{-1 ; p}}{p}\right)\right)-4\left(\left(\frac{p^{-1 ; q}}{q}\right)\right) & =\left(-4 \frac{q^{-1 ; p}}{p}+2\right)+\left(4 \frac{a^{-1 ; q}}{q}-2\right) \\
& =-4 \frac{q^{-1 ; p}}{p}+4 \frac{a^{-1 ; q}}{q}=4 \frac{-q q^{-1 ; p}+a^{-1 ; q} p}{p q} . \tag{5.2}
\end{align*}
$$

Then, using the equality $q^{-1 ; p} q=1+a^{-1 ; q} p$ (see Proposition 7.1), we obtain

$$
\begin{align*}
&-4\left(\left(\frac{q^{-1 ; p}}{p}\right)\right)-4\left(\left(\frac{p^{-1 ; q}}{q}\right)\right)=4 \frac{-q q^{-1 ; p}+a^{-1 ; q} p}{p q} \\
&=4 \frac{-(1+\alpha p)+a^{-1 ; q} p}{p q}=-\frac{4}{p q} \tag{5.3}
\end{align*}
$$

Finally,

$$
\begin{equation*}
-4\left(\left(\frac{q}{p}\right)\right)-4\left(\left(\frac{p}{q}\right)\right)=\left(-4 \frac{q}{p}+2\right)+\left(4 \frac{a}{q}-2\right)=-4 \frac{q}{p}+4 \frac{a}{q} . \tag{5.4}
\end{equation*}
$$

Next, we will prove useful reciprocity formulas for $N(q, p)$. Denote

$$
\begin{align*}
& R^{+}(q, p)=N(q, p)+N(p, q) \\
& R^{-}(q, p)=N(-q, p)+N(-p, q) \tag{5.5}
\end{align*}
$$

Proposition 5.2. Writing $p=e q-a$ with $0<a<q$, we have the following formulas:

$$
\begin{align*}
& R^{+}(q, p)= \begin{cases}-4 & \text { when } q=1 \text { and } p=2 \\
-14 & \text { when } 1<q=p-1 \\
4 p-14 & \text { when } 1=q<p-1 \\
4 e-22 & \text { when } p=e q-1 \\
4 e-24 & \text { when } 2 \leq a \leq q-1\end{cases}  \tag{5.6}\\
& R^{-}(q, p)= \begin{cases}-4 & \text { when } q=1 \text { and } p=2 \\
-6 & \text { when } 1<q=p-1, \\
-4 p+4 & \text { when } 1=q<p-1, \\
-4 e+2 & \text { when } p=e q-(q-1) \text { and } 1<q<p-1 \\
-4 e & \text { when } 1 \leq a \leq q-2 \text { and } 2<q\end{cases} \tag{5.7}
\end{align*}
$$

Proof. The first three formulas for both $R^{+}(q, p)$ and $R^{-}(q, p)$ are easily computable from the cases where $N(q, p)$ is easy to compute. Denote by $C_{(\alpha, \beta)}$ the constant term in $N(\alpha, \beta)$, so

$$
C_{(\alpha, \beta)}= \begin{cases}-6 & \text { for a non-exceptional singularity }  \tag{5.8}\\ -4 & \text { for an exceptional singularity }\end{cases}
$$

When $p=e q-a$, where $1 \leq a<q-1$, we have

$$
\begin{aligned}
R^{+}(q, p)=N(q, p)+N(p, q)= & {\left[C_{(q, p)}+A(q, p)-4\left(\left(\frac{q^{-1 ; p}}{p}\right)\right)-4\left(\left(\frac{q}{p}\right)\right)\right] } \\
& +\left[C_{(p, q)}+A(p, q)-4\left(\left(\frac{p^{-1 ; q}}{q}\right)\right)-4\left(\left(\frac{p}{q}\right)\right)\right]
\end{aligned}
$$

Then, by Proposition 5.1, we see that

$$
\begin{aligned}
R^{+}(q, p) & =C_{(q, p)}+C_{(p, q)}+\left[-12+4 e-4 \frac{a}{q}+4 \frac{q}{p}+\frac{4}{p q}\right]+\left[-\frac{4}{p q}-4 \frac{q}{p}+4 \frac{a}{q}\right] \\
& =4 e+C_{(q, p)}+C_{(p, q)}-12
\end{aligned}
$$

which proves the reciprocity formulas in each respective case. The proof for $R^{-}(q, p)$ is similar and is omitted.

We next use the above reciprocity relations to recursively compute an explicit formula for $N(q, p)$ :
Theorem 5.3. For $q$ and $p$ and relatively prime, we have

$$
N(q, p)= \begin{cases}\sum_{i=1}^{k} 4 e_{i}-12 k-2 & \text { when } q \not \equiv(p-1) \bmod p  \tag{5.9}\\ \sum_{i=1}^{k} 4 e_{i}-12 k=-4 p+4 & \text { when } q \equiv(p-1) \bmod p\end{cases}
$$

where $k$ and $e_{i}, 1 \leq i \leq k$, were defined above in the modified Euclidean algorithm (1.8).
Proof. We have already proved the second case in (4.3), and we will now prove the first case, so we need only consider $q \not \equiv(p-1) \bmod p$. Since our formulas only depend upon $q \bmod p$, we can assume that $1 \leq q<p-1$. We begin by using Proposition 5.2 to compute $N(q, p)$ as follows:

$$
\begin{aligned}
N(q, p) & =R^{+}(q, p)-N(p, q)=R^{+}(q, p)-N\left(e_{1} q-a_{1}, q\right) \\
& =R^{+}(q, p)-N\left(-a_{1}, q\right) \\
& =R^{+}(q, p)-N\left(-a_{1}, q\right)-N\left(-q, a_{1}\right)+N\left(-q, a_{1}\right) \\
& =R^{+}(q, p)-R^{-}\left(a_{1}, q\right)+N\left(a_{2}, a_{1}\right)+N\left(a_{1}, a_{2}\right)-N\left(a_{1}, a_{2}\right) \\
& =R^{+}(q, p)-R^{-}\left(a_{1}, q\right)+R^{+}\left(a_{2}, a_{1}\right)-N\left(-a_{3}, a_{2}\right) .
\end{aligned}
$$

Continuing this iteratively, we arrive at the formula

$$
\begin{aligned}
N(q, p)= & \sum_{i=1}^{r+1} 4 e_{i}-24\left\lceil\frac{r+1}{2}\right\rceil \\
& +\left[(-1)^{r+1} R^{(-1)^{r+1}}\left(a_{r+1}, a_{r}\right)+(-1)^{r+2} N\left((-1)^{r+2} a_{r+2}, a_{r+1}\right)\right]
\end{aligned}
$$

where $a_{r}=e_{r+2} a_{r+1}-1$ or $a_{r}=e_{r+2} a_{r+1}-\left(a_{r+1}-1\right)$. It is only necessary to consider the following four cases:
(1) When $r+2$ is even and $a_{r+2}=1$,

$$
N(q, p)=\sum_{i=1}^{r+3} 4 e_{i}-12(r+2)-14
$$

(2) When $r+2$ is odd and $a_{r+2}=1$,

$$
N(q, p)=\sum_{i=1}^{r+3} 4 e_{i}-12(r+1)-26
$$

(3) When $r+2$ is even and $a_{r+2}=a_{r+1}-1$,

$$
N(q, p)=\sum_{i=1}^{r+2} 4 e_{i}-4 a_{r+1}-12(r+2)+2 .
$$

(4) When $r+2$ is odd and $a_{r+2}=a_{r+1}-1$,

$$
N(q, p)=\sum_{i=1}^{r+2} 4 e_{i}-4 a_{r+1}-12(r+1)-10
$$

The formulas for $N(q, p)$ in each case are a direct consequence of (4.3) and Proposition 5.2. In cases (1) and (2), $k=r+3$. So in terms of $k$ we have

$$
\begin{equation*}
N(q, p)=\sum_{i=1}^{k} 4 e_{i}-12 k-2 \tag{5.10}
\end{equation*}
$$

for both cases. Now, in case (3), $k=\left(a_{r+1}-1\right)+(r+2)$ and $e_{i}=2$ for $i \geq r+3$. Therefore we can check that

$$
\begin{align*}
\sum_{i=1}^{k} 4 e_{i}-12 k-2 & =\left[\sum_{i=1}^{r+2} 4 e_{i}-12(r+2)\right]+\left[\sum_{i=r+3}^{k} 4 e_{i}-12\left(a_{r+1}-1\right)-2\right] \\
& =\sum_{i=1}^{r+2} 4 e_{i}-12(r+2)-4 a_{r+1}+2=N(q, p) \tag{5.11}
\end{align*}
$$

Finally, in case (4), $k=\left(a_{r+1}-1\right)+(r+2)$ and $e_{i}=2$ for $i \geq r+3$, and the result holds similarly.
Theorem 1.4 is then a trivial consequence of Theorem 5.3 and (4.1).

Remark 5.4. Ashikaga and Ishizaka [AI08, Theorem 1.1] prove a recursive formula for the Dedekind sum, which is equivalent to Theorem 5.3. However, our proof is more elementary and relies only on the reciprocity law for Dedekind sums. We will also need to use Proposition 5.2 below in Section 7.

## 6. Index on Calderbank-Singer spaces

In this section, we prove the results regarding the Calderbank-Singer metrics. Let $k$ and $k^{\prime}$ be the lengths of the modified Euclidean algorithm for $(q, p)$ and $(p-q, p)$ respectively.
Proof of Theorem 1.8. It follows from (1.12) that the compactified Calderbank-Singer space $(\hat{X}, \hat{g})$ satisfies $\tau_{\text {top }}(\hat{X})=-k$ and $\chi_{\text {top }}(\hat{X})=k+2$, so for a $(p-q, p)$-action when $q \neq 1$, the index is

$$
\operatorname{Ind}(\hat{X}, \hat{g})=\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)+N(q, p)=[-7 k+15]+\left[\sum_{i=1}^{k^{\prime}} 4 e_{i}^{\prime}-12 k^{\prime}-2\right]
$$

We next use a 4-dimensional $(q, p)$-football, denoted by $S_{(q, p)}^{4}$, to relate $k$ and $k^{\prime}$. This is defined using the $\Gamma_{(p, q)}$ action, acting as rotations around the $x_{5}$-axis:

$$
\begin{equation*}
S_{(q, p)}^{4}=S^{4} / \Gamma_{(q, p)} . \tag{6.1}
\end{equation*}
$$

This quotient is an orbifold with two singular points, one of ( $q, p$ )-type, and the other of $(-q, p)$-type. Since $\chi_{\text {top }}\left(S_{(q, p)}^{4}\right)=2$ and $\tau_{\text {top }}\left(S_{(q, p)}^{4}\right)=0$, the index of (1.1) on $S_{(q, p)}^{4}$ with the round metric $g_{S}$ is

$$
\begin{equation*}
\operatorname{Ind}\left(S_{(q, p)}^{4}, g_{S}\right)=3 \quad \text { for } 1<q<p-1 \tag{6.2}
\end{equation*}
$$

This is immediate from examining formula (3.5) for the index because the trigonometric sums coming from the ( $q, p$ )-type singularity cancel those coming from the ( $-q, p$ )-type singularity. Alternatively, this index can be computed directly, since it is well known that $H^{1}=0$ and $H^{2}=0$, and the dimension of the conformal automorphsim group of $S^{4} / \Gamma$ is 3 (for $1<q<p-1$ ), which follows from [McC02]. Using the formula

$$
\begin{equation*}
\operatorname{Ind}\left(S_{(q, p)}^{4}, g_{S}\right)=\frac{1}{2}\left(15 \chi_{\mathrm{top}}+29 \tau_{\mathrm{top}}\right)+N(q, p)+N(-q, p) \tag{6.3}
\end{equation*}
$$

and Theorem 5.3, we have

$$
\begin{align*}
-12 & =N(q, p)+N(-q, p)=N(q, p)+N(p-q, p) \\
& =\left[\sum_{i=1}^{k} 4 e_{i}-12 k-2\right]+\left[\sum_{i=1}^{k^{\prime}} 4 e_{i}^{\prime}-12 k^{\prime}-2\right], \tag{6.4}
\end{align*}
$$

which yields the formula

$$
\begin{equation*}
k^{\prime}=\frac{1}{12}\left(8+\sum_{i=1}^{k} 4 e_{i}+\sum_{i=1}^{k^{\prime}} 4 e_{i}^{\prime}-12 k\right) . \tag{6.5}
\end{equation*}
$$

Then, substituting this for $k$ in $\operatorname{Ind}(\hat{X}, \hat{g})$ gives

$$
\begin{align*}
\operatorname{Ind}(\hat{X}, \hat{g}) & =[-7 k+15]+\left[\sum_{i=1}^{k^{\prime}} 4 e_{i}^{\prime}-\left(8+\sum_{i=1}^{k} 4 e_{i}+\sum_{i=1}^{k^{\prime}} 4 e_{i}^{\prime}-12 k\right)-2\right] \\
& =5 k+5-\sum_{i=1}^{k} 4 e_{i} \tag{6.6}
\end{align*}
$$

Next, when $q=1$, we have $k=1$, so the index is

$$
\begin{equation*}
\operatorname{Ind}(\hat{X}, \hat{g})=[-7 k+15]+[-4 p+4]=-4 p+12 \tag{6.7}
\end{equation*}
$$

Proof of Theorem 1.9. Calderbank-Singer showed that their toric metrics come in families of dimension $k-1$. It was proved by Dominic Wright [Wri11, Corollary 1.1] that the moduli space of toric anti-self-dual metrics on the orbifolds is of dimension exactly $k-1$. So as long as we show the moduli space has dimension strictly larger than $k-1$, there must be non-toric deformations.

The $(1,2)$ case is the Eguchi-Hanson metric which has no deformations. For $q=1$ and $p>2$, the $(1, p)$-type Calderbank-Singer spaces are exactly the LeBrun negative mass metrics on $\mathcal{O}(-p)$ found in [LeB88]. For $p=3$, it was shown in [Hon13] that the moduli space of these metrics is of dimension 1 , so the result is true since $1>0=k-1$. For $p \geq 4$, by [Via13, Theorem 1.9], the moduli space has dimension at least $4 p-12>0$ (in fact the dimension is exactly $4 p-12$, see [Hon13, Theorem 1.1]). So the result holds for $q=1$ and $p \geq 3$. We also mention that [Hon13, Theorem 1.1] determines exactly the identity component of the automorphism groups of the deformations.

Next, assume that $q=p-1$. In this case, the metrics are hyperkähler, and correspond to toric multi-Eguchi-Hanson metrics. In this case, the moduli space of all hyperkähler metrics is known to be exactly of dimension $3(k-1)$.

Next, we assume that $1<q<p-1$. As mentioned in the Introduction, from [LM08, Theorem 4.2], we know that $\operatorname{dim}\left(H^{2}(\hat{X}, \hat{g})\right)=0$. Also, $\operatorname{dim}\left(H^{0}\right)=2$, since the metrics are toric and $q>1$. Therefore

$$
\begin{equation*}
\operatorname{dim}\left(H^{1}\right)=-\operatorname{Ind}(\hat{X}, \hat{g})+\operatorname{dim}\left(H^{0}\right)=-\operatorname{Ind}(\hat{X}, \hat{g})+2 . \tag{6.8}
\end{equation*}
$$

When $q \neq 1$, we have

$$
\begin{equation*}
-\mathrm{Ind}=-5 k-5+\sum_{i=1}^{k} 4 e_{i} \tag{6.9}
\end{equation*}
$$

Since $e_{i} \geq 2$ for all $i$ and since $q<p-1$, we have $e_{j} \geq 3$ for some $j, 1 \leq j \leq k$. Therefore

$$
\begin{equation*}
\operatorname{dim}\left(H^{1}\right) \geq 3 k+1 \tag{6.10}
\end{equation*}
$$

The actual moduli space is locally isomorphic to $H^{1} / H^{0}$, so it has dimension at least $3 k-1>3(k-1)$.

## 7. Index on weighted projective spaces

In this section we will study the index of the complex (1.1) at the Bochner-Kähler metrics of Bryant with reversed orientation to make them anti-self-dual. This reversal of orientation makes the orbifold points have orientation-reversing conjugate actions as follows:
(1) Around $[1,0,0]$ there is a $\left(-q^{-1 ; r} p, r\right)$-action.
(2) Around $[0,1,0]$ there is a $\left(-p^{-1 ; q} r, q\right)$-action.
(3) Around $[0,0,1]$ there is a $\left(-r^{-1 ; p} q, p\right)$-action.

In the next subsection, we will present some elementary number-theoretic propositions that we will use throughout our computations. After that, we will prove crucial reciprocity laws for sawtooth functions relating $r, q$ and $p$ and then employ these to prove our main formula for the index. Finally, we use this formula to prove Theorem 1.13.

### 7.1. Elementary number-theoretic preliminaries

Recall that for two relatively prime positive integers $1<\alpha<\beta$, we denote $\alpha$ 's inverse modulo $\beta$ by $\alpha^{-1 ; \beta}$, and $\beta$ 's inverse modulo $\alpha$ by $\beta^{-1 ; \alpha}$. Since $\alpha<\beta$ we can write

$$
\begin{equation*}
\beta=e \alpha-a \tag{7.1}
\end{equation*}
$$

where $e$ and $a$ are positive integers with $a<\alpha$. Then we have the following proposition:
Proposition 7.1. We have the following identities:
(1) $\beta^{-1 ; \alpha}=\alpha-a^{-1 ; \alpha}$,
(2) $\alpha \alpha^{-1 ; \beta}=1+a^{-1 ; \alpha} \beta$.

Proof. To prove the first identity, recall that $\beta=e \alpha-a$, so

$$
\beta\left(\alpha-a^{-1 ; \alpha}\right)=(e \alpha-a)\left(\alpha-a^{-1 ; \alpha}\right)=e \alpha^{2}-e \alpha a^{-1 ; \alpha}-a \alpha+a a^{-1 ; \alpha} \equiv 1 \bmod \alpha .
$$

This proves the first identity because $\alpha-a^{-1 ; \alpha}<\alpha$ and the multiplicative inverses are unique.

To prove the second identity we first write

$$
\begin{equation*}
\alpha \alpha^{-1 ; \beta}=1+X \beta . \tag{7.2}
\end{equation*}
$$

Since $\alpha>1$ we know $X$ must be a positive integer. We can then solve for

$$
\begin{equation*}
\beta=\frac{\alpha \alpha^{-1 ; \beta}-1}{X} . \tag{7.3}
\end{equation*}
$$

Therefore $\beta=\frac{\alpha \alpha^{-1 ; \beta}-1}{X}=e \alpha-a$, so

$$
\begin{equation*}
\alpha \alpha^{-1 ; \beta}-1=e \alpha X-a X, \tag{7.4}
\end{equation*}
$$

from which we see that

$$
\begin{equation*}
a X=\alpha\left(e X-\alpha^{-1 ; \beta}\right)+1, \tag{7.5}
\end{equation*}
$$

so $a X \equiv 1 \bmod \alpha$. This proves the second identity because $X=\frac{\alpha \alpha^{-1 ; \beta}-1}{\beta}<\alpha$ and multiplicative inverses are unique.

Recall that the fractional part of $x$ is defined by

$$
\begin{equation*}
\{x\}=x-\lfloor x\rfloor . \tag{7.6}
\end{equation*}
$$

We will use the following proposition extensively in the next section, the proof is elementary:

Proposition 7.2. For any real $\alpha$ and $\beta$, both non-integral,

$$
((\alpha+\beta))= \begin{cases}((\alpha))+((\beta))+1 / 2 & \text { when }\{\alpha\}+\{\beta\}<1  \tag{7.7}\\ ((\alpha))+((\beta))-1 / 2 & \text { when }\{\alpha\}+\{\beta\}>1 \\ 0 & \text { when }\{\alpha\}+\{\beta\}=1\end{cases}
$$

### 7.2. Reciprocity formulas for sawtooth functions

Let $r<q<p$ and write

$$
\begin{align*}
& p=e_{p r} r-a_{p r} \\
& p=e_{p q} q-a_{p q}  \tag{7.8}\\
& q=e_{q r} r-a_{q r}
\end{align*}
$$

We have the following identities from Proposition 7.1(1):

$$
\begin{align*}
& p^{-1 ; r}=r-a_{p r}^{-1 ; r} \\
& p^{-1 ; q}=q-a_{p q}^{-1 ; q}  \tag{7.9}\\
& q^{-1 ; r}=r-a_{q r}^{-1 ; r}
\end{align*}
$$

and from Proposition 7.1(2):

$$
\begin{align*}
& r r^{-1 ; p}=1+a_{p r}^{-1 ; r} p \\
& r r^{-1 ; q}=1+a_{q r}^{-1 ; r} q  \tag{7.10}\\
& q q^{-1 ; p}=1+a_{p q}^{-1 ; q} p
\end{align*}
$$

We now use these identities to prove reciprocity laws for the sawtooth function. These reciprocity laws will be broken up into two theorems where the first is independent of $r+q$ in relation to $p$ and the second is dependent.

Theorem 7.3. We have the following reciprocity relations:
(1) $\left(\left(\frac{q p^{-1 ; r}}{r}\right)\right)+\left(\left(\frac{q r^{-1 ; p}}{p}\right)\right)=\frac{q}{p r}$,
(2) $\left(\left(\frac{r p^{-1 ; q}}{q}\right)\right)+\left(\left(\frac{r q^{-1 ; p}}{p}\right)\right)=\frac{r}{p q}$.

Proof. Consider the first reciprocity relation. We have

$$
\begin{align*}
\left(\left(\frac{q p^{-1 ; r}}{r}\right)\right)+\left(\left(\frac{q r^{-1 ; p}}{p}\right)\right) & =\left(\left(\frac{q\left(r-a_{p r}^{-1 ; r}\right)}{r}\right)\right)+\left(\left(\frac{r r^{-1 ; p} q}{p r}\right)\right) \\
& =\left(\left(q-\frac{q a_{p r}^{-1 ; r}}{r}\right)\right)+\left(\left(\frac{q}{p r}+\frac{q a_{p r}^{-1 ; r}}{r}\right)\right) \\
& =-\left(\left(\frac{q a_{p r}^{-1 ; r}}{r}\right)\right)+\left(\left(\frac{q}{p r}+\frac{q a_{p r}^{-1 ; r}}{r}\right)\right) \tag{7.11}
\end{align*}
$$

Now, we can write $q a_{p r}^{-1 ; r} / r=X+C / r$ where $X>0$ and $0<C<r$ are positive integers so that

$$
\begin{equation*}
-\left(\left(\frac{q a_{p r}^{-1 ; r}}{r}\right)\right)+\left(\left(\frac{q}{p r}+\frac{q a_{p r}^{-1 ; r}}{r}\right)\right)=-\left(\left(\frac{C}{r}\right)\right)+\left(\left(\frac{q}{p r}+\frac{C}{r}\right)\right) \tag{7.12}
\end{equation*}
$$

Since $0<C<r$ we know that

$$
\begin{equation*}
\frac{q}{p r}+\frac{C}{r} \leq \frac{q}{p r}+\frac{r-1}{r}=\frac{q}{p r}+\frac{p r-p}{p r}<1 \tag{7.13}
\end{equation*}
$$

because $q<p$, which implies that

$$
\begin{equation*}
\left\{\frac{q}{p r}\right\}+\left\{\frac{C}{r}\right\}<1 \tag{7.14}
\end{equation*}
$$

Therefore, by Proposition 7.2, we can separate the second sawtooth function to get

$$
\begin{equation*}
\left(\left(\frac{q}{p r}+\frac{C}{r}\right)\right)=\left(\left(\frac{q}{p r}\right)\right)+\left(\left(\frac{C}{r}\right)\right)+\frac{1}{2} \tag{7.15}
\end{equation*}
$$

Putting this back into (7.12) we see that

$$
\begin{align*}
\left(\left(\frac{q p^{-1 ; r}}{r}\right)\right) & +\left(\left(\frac{q r^{-1 ; p}}{p}\right)\right)=-\left(\left(\frac{C}{r}\right)\right)+\left(\left(\frac{q}{p r}+\frac{C}{r}\right)\right) \\
& \left.=-\left(\left(\frac{C}{r}\right)\right)+\left(\left(\frac{q}{p r}\right)\right)+\left(\left(\frac{C}{r}\right)\right)\right)+\frac{1}{2}=\left(\left(\frac{q}{p r}\right)\right)+\frac{1}{2} \\
& =\frac{q}{p r}-\left\lfloor\frac{q}{p r}\right\rfloor=\frac{q}{p r} \tag{7.16}
\end{align*}
$$

The proof of the second reciprocity relation exactly follows the proof of the first.
The next theorem gives a similar reciprocity relation, but it is dependent upon $r+q$ in relation to $p$.

Theorem 7.4. For $1<r<q<p$ we have the following reciprocity relation:

$$
\left(\left(\frac{q^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{r^{-1 ; q} p}{q}\right)\right)= \begin{cases}\frac{p}{q r} & \text { when } r+q>p  \tag{7.17}\\ \frac{p}{q r}-1 & \text { when } r+q=p \\ \frac{p}{q r}-\left\lfloor\frac{p}{q r}\right\rfloor & \text { when } r+q<p \\ & \text { and }\left\{\frac{p}{q r}\right\}<\left\{\frac{q^{-1 ; r} p}{r}\right\} \\ \frac{p}{q r}-\left\lfloor\frac{p}{q r}\right\rfloor-1 & \text { when } r+q<p \\ & \text { and }\left\{\frac{p}{q r}\right\}>\left\{\frac{q^{-1 ; r} p}{r}\right\}\end{cases}
$$

Proof. We begin in a similar way to the proof of Theorem 7.3:

$$
\begin{align*}
\left(\left(\frac{q^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{r^{-1 ; q} p}{q}\right)\right) & =\left(\left(\frac{p\left(r-a_{q r}^{-1 ; r}\right)}{r}\right)\right)+\left(\left(\frac{\left(1+a_{q r}^{-1 ; r} q\right) p}{q r}\right)\right) \\
& =-\left(\left(\frac{a_{q r}^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{p}{q r}+\frac{a_{q r}^{-1 ; r} p}{r}\right)\right) \tag{7.18}
\end{align*}
$$

Now, we can write $a_{q r}^{-1 ; r} p / r=X+C / r$ where $X>0$ and $0<C<r$ are positive integers so that

$$
\begin{equation*}
-\left(\left(\frac{a_{q r}^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{p}{q r}+\frac{a_{q r}^{-1 ; r} p}{r}\right)\right)=-\left(\left(\frac{C}{r}\right)\right)+\left(\left(\frac{p}{q r}+\frac{C}{r}\right)\right) \tag{7.19}
\end{equation*}
$$

The same argument that we used in the previous proof to split up the second sawtooth function will no longer work because $p>q$, which could allow $p /(q r)+C / r>1$. Let us consider the first case of this reciprocity relation when $r+q>p$. In this case we know that $p<r q$ so $p<2 q$ since $1<r<q$. Now, we will show that $C \leq r-2$ and use this to prove the first case. Write $p$ as

$$
\begin{equation*}
p=k_{p r} r+m_{p r} \tag{7.20}
\end{equation*}
$$

where $k_{p r}=e_{p r}-1$ and $m_{p r}=r-a_{p r}$ are positive integers. We know that $C / r$ is going to be the fractional part of $p a_{q r}^{-1 ; r} / r$ which equals the fractional part of $m_{p r} a_{q r}^{-1 ; r} / r$. If this equals $(r-1) / r$ then $m_{p r} a_{p r}^{-1 ; r} \equiv-1 \bmod r$ and therefore $m_{p r}=r-a_{q r}$, because multiplicative inverses are unique, which implies that $a_{p r}=a_{q r}$ (denote this value by $A$ ). So we have

$$
\begin{equation*}
e_{p r} r-A=p<q+r=\left(e_{q} r-A\right)+r \tag{7.21}
\end{equation*}
$$

which is a contradiction because $e_{p r} \geq e_{q r}+1$. Therefore $C / r \leq(r-2) / r$, so

$$
\begin{equation*}
\frac{p}{q r}+\frac{C}{r} \leq \frac{p}{q r}+\frac{r-2}{r}=\frac{p}{q r}+\frac{q r-2 q}{q r}<1 \tag{7.22}
\end{equation*}
$$

which implies that

$$
\begin{equation*}
\left\{\frac{p}{q r}\right\}+\left\{\frac{C}{r}\right\}<1 \tag{7.23}
\end{equation*}
$$

Therefore, by Proposition 7.2, we can separate the second sawtooth function to get

$$
\begin{align*}
-\left(\left(\frac{C}{r}\right)\right)+\left(\left(\frac{p}{q r}+\frac{C}{r}\right)\right) & =-\left(\left(\frac{C}{r}\right)\right)+\left(\left(\frac{p}{q r}\right)\right)+\left(\left(\frac{C}{r}\right)\right)+\frac{1}{2} \\
& =\left(\left(\frac{p}{q r}\right)\right)+\frac{1}{2}=\frac{p}{q r}-\left\lfloor\frac{p}{q r}\right\rfloor=\frac{p}{q r} \tag{7.24}
\end{align*}
$$

which proves the first case.
Now, consider the second case when $r+q=p$. In this case we see that

$$
\begin{align*}
\left(\left(\frac{p}{q r}+\frac{p a_{q r}^{-1 ; r}}{r}\right)\right) & =\left(\left(\frac{p}{q r}+\frac{(r+q) a_{q r}^{-1 ; r}}{r}\right)\right) \\
& =\left(\left(\frac{p}{q r}-\frac{a_{q r} a_{q r}^{-1 ; r}}{r}\right)\right)=\left(\left(\frac{p}{q r}-\frac{1}{r}\right)\right) \\
& =\left(\frac{p}{q r}-\frac{1}{r}\right)-\left\lfloor\frac{p}{q r}-\frac{1}{r}\right\rfloor-\frac{1}{2}=\frac{1}{q}-\frac{1}{2} \tag{7.25}
\end{align*}
$$

Then, we compare this to

$$
\begin{align*}
\left(\left(\frac{p}{q r}\right)\right)+ & \left(\left(\frac{p a_{q r}^{-1 ; r}}{r}\right)\right)+\frac{1}{2}=\left(\left(\frac{p}{q r}\right)\right)+\left(\left(\frac{(r+q) a_{q r}^{-1 ; r}}{r}\right)\right)+\frac{1}{2} \\
& =\left(\left(\frac{p}{q r}\right)\right)+\left(\left(\frac{-a_{q r} a_{q r}^{-1 ; r}}{r}\right)\right)+\frac{1}{2}=\left(\left(\frac{p}{q r}\right)\right)-\left(\left(\frac{1}{r}\right)\right)+\frac{1}{2} \\
& =\left(\frac{p}{q r}-\left\lfloor\frac{p}{q r}\right\rfloor-\frac{1}{2}\right)-\left(\frac{1}{r}-\left\lfloor\frac{1}{r}\right\rfloor-\frac{1}{2}\right)+\frac{1}{2} \\
& =\frac{p}{q r}-\frac{1}{r}+\frac{1}{2}=\frac{1}{q}+\frac{1}{2}=\left(\left(\frac{p}{q r}+\frac{p a_{q r}^{-1 ; r}}{r}\right)\right)+1 \tag{7.26}
\end{align*}
$$

Therefore we see that

$$
\begin{align*}
-\left(\left(\frac{a_{q r}^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{p}{q r}\right.\right. & \left.\left.+\frac{a_{q r}^{-1 ; r} p}{r}\right)\right) \\
& =-\left(\left(\frac{a_{q r}^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{p}{q r}\right)\right)+\left(\left(\frac{p a_{q r}^{-1 ; r}}{r}\right)\right)-\frac{1}{2} \\
& =\left(\left(\frac{p}{q r}\right)\right)-\frac{1}{2}=\frac{p}{q r}-\left\lfloor\frac{p}{q r}\right\rfloor-1=\frac{p}{q r}-1 \tag{7.27}
\end{align*}
$$

which proves the second case.

To prove the third and fourth cases we begin once again by using

$$
\begin{equation*}
\left(\left(\frac{q^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{r^{-1 ; q} p}{q}\right)\right)=\left(\left(\frac{q^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{p}{q r}-\frac{q^{-1 ; r} p}{r}\right)\right) \tag{7.28}
\end{equation*}
$$

Notice that

$$
\begin{equation*}
\left\{\frac{p}{q r}\right\}+\left\{\frac{-q^{-1 ; r} p}{r}\right\}=\left\{\frac{p}{q r}\right\}+1-\left\{\frac{q^{-1 ; r} p}{r}\right\} \tag{7.29}
\end{equation*}
$$

which is never equal to one because $r, q$ and $p$ are relatively prime. Now, the rest of the proof of the third and fourth cases follows directly from Proposition 7.2.

## 7.3. $\Gamma$-index for weighted projective spaces

First, recall Definition 4.3: singularities resulting from a $(p-1, p)$-action are said to be exceptional and otherwise they are called non-exceptional. Consider the case when $1<r<q<p$ so that there are three singularities. Before giving theorems concerning the index, we will first examine what type singularities, non-exceptional or exceptional, are admitted around each orbifold point in the cases when $r+q>p, r+q=p$ and $r+q<p$.

Proposition 7.5. When $r+q>p$ all three singularities are non-exceptional. If $r+q=p$ then:
(1) The singularity at $[1,0,0]$ is always exceptional.
(2) The singularity at $[0,1,0]$ is always exceptional.
(3) The singularity at $[0,0,1]$ is non-exceptional and comes from a $(1, p)$-action.

If $r+q<p$ then:
(1) The singularity at $[1,0,0]$ is exceptional if and only if $p \equiv q \bmod r$.
(2) The singularity at $[0,1,0]$ is exceptional if and only if $p \equiv r \bmod q$.
(3) The singularity at $[0,0,1]$ is always non-exceptional.

Proof. At $[1,0,0]$ the $\left(-q^{-1 ; r} p, r\right)$-action is equivalent to a $\left(-a_{q r}^{-1 ; r} a_{p r}, r\right)$-action, and this is equivalent to an $(r-1, r)$-action if and only if $a_{p r}=a_{q r}$. If $r+q>p$, suppose that $a_{p r}=a_{q r}$; then

$$
\begin{equation*}
p=e_{p r} r-a_{q r} \quad \text { and } \quad q=e_{q r} r-a_{q r}, \tag{7.30}
\end{equation*}
$$

so $p<q+r=\left(e_{q r}+1\right) r-a_{q r}$, which is a contradiction because $e_{p r} \geq e_{q r}+1$. If $r+q=p$ we have

$$
\begin{equation*}
p=q+r=\left(e_{q r}+1\right) r-a_{q r}, \tag{7.31}
\end{equation*}
$$

so $a_{p r}=a_{q r}$ since $a_{q r}<r$. If $r+q<p$, then this happens if and only if $p \equiv q \bmod r$.

At $[0,1,0]$, by Remark 1.3, the $\left(-p^{-1 ; q} r, q\right)$-action is equivalent to a $\left(-r^{-1 ; q} p, q\right)$ action. This is equivalent to an $\left(r^{-1 ; q} a_{p q}, q\right)$-action, which is equivalent to a $(q-1, q)$ action if and only if $a_{p q} r^{-1 ; q} \equiv-1 \bmod q$, which would imply that $a_{p q}=q-r$. If $r+q>p$, suppose that $a_{p q}=q-r$; then

$$
\begin{equation*}
p=2 q-a_{p q}=2 q-(q-r)=q+r \tag{7.32}
\end{equation*}
$$

which is a contradiction because $r+q>p$. If $r+q=p$, we have

$$
\begin{equation*}
p=2 q-(q-r), \tag{7.33}
\end{equation*}
$$

so $a_{p q}=q-r$. If $r+q<p$ then this happens if and only if $p \equiv r \bmod q$.
At $[0,0,1]$ the $\left(-r^{-1 ; p} q, p\right)$-action is equivalent to a $(p-1, p)$-action if and only if $r^{-1 ; p} q \equiv 1 \bmod p$. If $r+q>p$, this condition would imply that $q=r$, which is a contradiction. If $r+q=p$ then the $\left(-r^{-1 ; p} q, p\right)$-action is obviously equivalent to a $(1, p)$-action since $q=p-r$. If $r+q<p$ then $r^{-1 ; p} q \equiv 1 \bmod p$ occurs if and only if $q=r$, but $q>r$ so this can never happen.
In the case $r+q<p$, we can add the following:
Proposition 7.6. When $r+q<p$ and the singularities at $[1,0,0]$ and $[0,1,0]$ are both exceptional, we have $p=X q r+r+q$ for some integer $X$, and

$$
\begin{equation*}
\left\{\frac{p}{q r}\right\}>\left\{\frac{q^{-1 ; r} p}{r}\right\} \tag{7.34}
\end{equation*}
$$

Proof. Since the singularities around $[1,0,0]$ and $[0,1,0]$ are both exceptional, from Proposition 7.5 we know that

$$
\begin{equation*}
p \equiv q \bmod r \quad p \equiv r \bmod q \tag{7.35}
\end{equation*}
$$

Therefore, we can write

$$
\begin{equation*}
p=Y_{1} q+r=Y_{2} r+q, \tag{7.36}
\end{equation*}
$$

and solve for

$$
\begin{equation*}
r=\frac{Y_{1}-1}{Y_{2}-1} q, \tag{7.37}
\end{equation*}
$$

which implies that $q X=Y_{2}-1$ for some $X$ in $\mathbb{Z}$, since $q$ and $r$ are relatively prime. Then solving for $Y_{2}=q X+1$ we see that

$$
\begin{equation*}
p=(q X+1) r+q=X q r+r+q . \tag{7.38}
\end{equation*}
$$

Now, since $p=X q r+r+q$ we see that $a_{p r}=a_{q r}$. Therefore

$$
\begin{align*}
\left\{\frac{p}{q r}\right\}-\left\{\frac{a_{q r}^{-1 ; r} a_{p r}}{r}\right\} & =\left\{\frac{X q r+r+q}{q r}\right\}-\left\{\frac{a_{q r}^{-1 ; r} a_{q r}}{r}\right\} \\
& =\left\{\frac{1}{q}+\frac{1}{r}\right\}-\left\{\frac{1}{r}\right\}=\frac{1}{q}>0 \tag{7.39}
\end{align*}
$$

The following is the main result of this section, which is the same as Theorem 1.12 upon identifying the integer $\epsilon$ with the number of exceptional singularities:

Theorem 7.7. Let $g$ be the canonical Bochner-Kähler metric with reversed orientation on $\overline{\mathbb{C P}}_{(r, q, p)}^{2}$, and assume that $1<r<q<p$. If $r+q \geq p$ then

$$
\begin{equation*}
\operatorname{Ind}\left(\overline{\mathbb{C P}}_{(r, q, p)}^{2}, g\right)=2 \tag{7.40}
\end{equation*}
$$

If $r+q<p$ then

$$
\operatorname{Ind}\left(\overline{\mathbb{C P}}_{(r, q, p)}^{2}, g\right)= \begin{cases}2+2 \epsilon-4\left\lfloor\frac{p}{q r}\right\rfloor & \text { when }\left\{\frac{p}{q r}\right\}<\left\{\frac{q^{-1 ; r} p}{r}\right\},  \tag{7.41}\\ -2+2 \epsilon-4\left\lfloor\frac{p}{q r}\right\rfloor & \text { when }\left\{\frac{p}{q r}\right\}>\left\{\frac{q^{-1 ; r}}{r}\right\},\end{cases}
$$

where $\epsilon$ is the number of exceptional singularities, either 0,1 , or 2 .
Note that from Proposition 7.6 the only instance when two exceptional singularities can occur is in the second case, thus there are really only five distinct cases. All of these cases do in fact occur: see Table 7.1.

Table 7.1. Cases in Theorem 7.7

| $(r, q, p)$ | $\epsilon$ | $\left\{\frac{p}{q r}\right\}$ |
| :---: | :---: | :---: |
| $(3,7,11)$ | 0 | $\left.<0 \frac{q^{-1 ; r}}{r}\right\}$ |
| $(3,7,41)$ | 0 | $>0$ |
| $(3,7,25)$ | 1 | $<0$ |
| $(3,7,13)$ | 1 | $>0$ |
| $(3,7,31)$ | 2 | $>0$ |

Proof of Theorem 7.7. Since $1<r<q<p$, there are three singularities. Furthermore, $\chi_{\text {top }}=3$ and $\tau_{\text {top }}=-1($ see $[\operatorname{Dim} 92$, Appendix B]), so the $\Gamma$-index is

$$
\begin{align*}
\text { Ind }= & 8+N\left(-q^{-1 ; r} p, r\right)+N\left(-p^{-1 ; q} r, q\right)+N\left(-r^{-1 ; p} q, p\right) \\
= & 8+\left[C_{\left(-q^{-1 ; r} p, r\right)}+A\left(-q^{-1 ; r} p, r\right)-4\left(\left(\frac{-q^{-1 ; r} p}{r}\right)\right)-4\left(\left(\frac{-p^{-1 ; r} q}{r}\right)\right)\right] \\
& +\left[C_{\left(-r^{-1 ; q} p, q\right)}+A\left(-p^{-1 ; q} r, q\right)-4\left(\left(\frac{-r^{-1 ; q} p}{q}\right)\right)-4\left(\left(\frac{-p^{-1 ; q} r}{q}\right)\right)\right] \\
& +\left[C_{\left(-r^{-1 ; p} q, p\right)}+A\left(-r^{-1 ; p} q, p\right)-4\left(\left(\frac{-r^{-1 ; p} q}{p}\right)\right)-4\left(\left(\frac{-q^{-1 ; p} r}{p}\right)\right)\right] \tag{7.42}
\end{align*}
$$

recalling that $C_{(\alpha, \beta)}$ was defined above in (5.8). Then, using Rademacher's triple reciprocity for Dedekind sums [Rad54]

$$
\begin{equation*}
s\left(q^{-1 ; r} p, r\right)+s\left(p^{-1 ; q} r, q\right)+s\left(r^{-1 ; p} q, p\right)=-\frac{1}{4}+\frac{1}{12}\left(\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}\right) \tag{7.43}
\end{equation*}
$$

we see that

$$
\begin{align*}
\text { Ind }= & 8+\left[C_{\left(-q^{-1 ; r} p, r\right)}+C_{\left(-r^{-1 ; q} p, q\right)}+C_{\left(-r^{-1 ; p} q, p\right)}\right] \\
& +48\left[\frac{1}{4}-\frac{1}{12}\left(\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}\right)\right] \\
& +4\left[\left(\left(\frac{q^{-1 ; r} p}{r}\right)\right)+\left(\left(\frac{p^{-1 ; r} q}{r}\right)\right)+\left(\left(\frac{r^{-1 ; q} p}{q}\right)\right)\right] \\
& +4\left[\left(\left(\frac{p^{-1 ; q} r}{q}\right)\right)+\left(\left(\frac{r^{-1 ; p} q}{p}\right)\right)+\left(\left(\frac{q^{-1 ; p} r}{p}\right)\right)\right] . \tag{7.44}
\end{align*}
$$

Now, using our reciprocity laws for sawtooth functions, Theorems 7.3 and 7.4, and the restrictions on the types of singularities admitted, Proposition 7.5, we complete the proof for each case.

When $r+q>p$,

$$
\text { Ind }=8+[-18]+48\left[\frac{1}{4}-\frac{1}{12}\left(\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}\right)\right]+4\left[\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}\right]=2 .
$$

When $r+q=p$,

$$
\text { Ind }=8+[-14]+48\left[\frac{1}{4}-\frac{1}{12}\left(\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}\right)\right]+4\left[\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}-1\right]=2
$$

When $r+q<p$ and $\left\{\frac{p}{q r}\right\}<\left\{\frac{q^{-1 ; r}}{r}\right\}$,

$$
\begin{aligned}
\text { Ind }= & 8+\left[C_{\left(-q^{-1 ; r} p, r\right)}+C_{\left(-r^{-1 ; q} p, q\right)}+C_{\left(-r^{-1 ; p} q, p\right)}\right] \\
& +48\left[\frac{1}{4}-\frac{1}{12}\left(\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}\right)\right]+4\left[\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}-\left\lfloor\frac{p}{q r}\right\rfloor\right] \\
= & 20+\left[C_{\left(-q^{-1 ; r} p, r\right)}+C_{\left(-r^{-1 ; q} p, q\right)}+C_{\left(-r^{-1 ; p} q, p\right)}\right]-4\left\lfloor\frac{p}{q r}\right\rfloor \\
= & 2+2 \epsilon-4\left\lfloor\frac{p}{q r}\right\rfloor .
\end{aligned}
$$

When $r+q<p$ and $\left\{\frac{p}{q r}\right\}>\left\{\frac{q^{-1 ; r}}{r}\right\}$,

$$
\begin{aligned}
\text { Ind }= & 8+\left[C_{\left(-q^{-1 ; r} p, r\right)}+C_{\left(-r^{-1 ; q} p, q\right)}+C_{\left(-r^{-1 ; p} q, p\right)}\right] \\
& +48\left[\frac{1}{4}-\frac{1}{12}\left(\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}\right)\right]+4\left[\frac{r}{p q}+\frac{q}{p r}+\frac{p}{q r}-1-\left\lfloor\frac{p}{q r}\right\rfloor\right] \\
= & 16+\left[C_{\left(-q^{-1 ; r} p, r\right)}+C_{\left(-r^{-1 ; q} p, q\right)}+C_{\left(-r^{-1 ; p} p, p\right)}\right]-4\left\lfloor\frac{p}{q r}\right\rfloor \\
= & -2+2 \epsilon-4\left\lfloor\frac{p}{q r}\right\rfloor .
\end{aligned}
$$

This completes the proof.

We also state the following theorem, which gives the index in the cases when there are strictly less than three singularities.

Theorem 7.8. Let $g$ be the canonical Bochner-Kähler metric with reversed orientation on $\overline{\mathbb{P}}_{(r, q, p)}^{2}$. When $1=r<q<p$ there are two singularities and

$$
\operatorname{Ind}\left(\overline{\mathbb{C P}}_{(1, q, p)}^{2}, g\right)= \begin{cases}2 & \text { when } q=p-1  \tag{7.45}\\ -4\left\lfloor\frac{p}{q}\right\rfloor+6 & \text { when } p=e q-(q-1) \text { and } q \neq p-1 \\ -4\left\lfloor\frac{p}{q}\right\rfloor+4 & \text { when } 1 \leq a_{p q} \leq q-2 \text { and } q>2\end{cases}
$$

When $1=r=q<p$ there is one singularity and

$$
\begin{equation*}
\operatorname{Ind}\left(\overline{\mathbb{C P}}_{(1,1, p)}^{2}, g\right)=-4 p+12 \tag{7.46}
\end{equation*}
$$

Proof. We have

$$
\begin{equation*}
\frac{1}{2}\left(15 \chi_{\text {top }}+20 \tau_{\text {top }}\right)=8, \tag{7.47}
\end{equation*}
$$

Since $1=r<q<p$ we know that $p>2$. The first case follows from the reciprocity formula for $R^{-}(q, p)$ in Proposition 5.2. The second case follows from $N(-1, p)=$ $-4 p+4$ in (4.3).

### 7.4. Proof of Theorem 1.13

We first present a general result about $H^{2}(M, g)$ on certain self-dual Kähler orbifolds:
Proposition 7.9. Let $(M, g)$ be a compact self-dual Kähler orbifold and assume that the set $M^{>0}=\{p \in M: R(p)>0\}$ is non-empty. With the reversed orientation to make $g$ anti-self-dual, we have $H^{2}(M, g)=0$.
Proof. As mentioned in the Introduction, the metric $\tilde{g}=R^{-2} g$ is an Einstein metric, which is complete on components of $M^{*}$. If $Z \in S_{0}^{2}\left(\Lambda_{+}^{2}\left(T^{*} M\right)\right)$ satisfies $\mathcal{D}_{g}^{*} Z=0$, where $\mathcal{D}_{g}^{*}$ is the adjoint of $\mathcal{D}_{g}$, then from conformal invariance $\mathcal{D}_{\tilde{g}}^{*} Z=0$ when $Z$ is viewed as a $(1,3)$ tensor. We compute

$$
\begin{equation*}
|Z|_{\tilde{g}}^{2}=\tilde{g}^{i p} \tilde{g}^{j q} Z_{i j k}^{l} Z_{p q l}{ }^{k}=R^{4} g^{i p} g^{j q} Z_{i j k}^{l} Z_{p q l}{ }^{k}=R^{4}|Z|_{g}^{2}, \tag{7.48}
\end{equation*}
$$

so we have

$$
\begin{equation*}
|Z|_{\tilde{g}}=R^{2}|Z|_{g} \tag{7.49}
\end{equation*}
$$

Let $M_{1}^{*}$ denote any non-trivial component of $M^{*}$. Since the metric $\tilde{g}$ is Einstein on $M_{1}^{*}$, from [Ito95, Proposition 5.1], we have

$$
\begin{equation*}
\mathcal{D}_{\tilde{g}} \mathcal{D}_{\tilde{g}}^{*} Z=\frac{1}{24}\left(3 \nabla_{\tilde{g}}^{*} \nabla_{\tilde{g}}+2 R_{\tilde{g}}\right)\left(2 \nabla_{\tilde{g}}^{*} \nabla_{\tilde{g}}+R_{\tilde{g}}\right) Z, \tag{7.50}
\end{equation*}
$$

where $R_{\tilde{g}}$ is the (constant) scalar curvature of the Einstein metric $\tilde{g}$ on $M_{1}^{*}$. If $M^{>0}=M$, then the maximum principle immediately implies that $Z=0$. Otherwise, there is a nontrivial open component of $M^{*}$, which we again call $M_{1}^{*}$. The metric $\tilde{g}$ is a complete Einstein metric on $M_{1}^{*}$, and (7.49) shows that $|Z|_{\tilde{g}}(x)=o(1)$ as $r \rightarrow 0$, where $r$ is the
distance to the zero set of the scalar curvature. Viewed on the complete manifold ( $M_{1}^{*}, \tilde{g}$ ), $Z$ is then a decaying solution at infinity of (7.50). Since $R_{\tilde{g}}$ is a constant, a standard separation of variables argument (see for example [Don89]) implies that $Z$ must decay faster than the inverse of any polynomial in the $\tilde{g}$ metric (it moreover has exponential decay). Equivalently, $|Z|_{g}=O\left(r^{k}\right)$ as $r \rightarrow 0$ for any $k>0$. This implies that $Z$ has a zero of infinite order along the zero set of the scalar curvature. The unique continuation principle for elliptic operators (see [Aro57]) then implies that $Z$ is identically zero.

Corollary 7.10. If $g$ is the canonical Bochner-Kähler metric with reversed orientation on $\overline{\mathbb{C P}}_{(r, q, p)}^{2}$, then $H^{2}(M, g)=0$.
Proof. From [DG06, equation (2.32)], the set $M^{>0}$ is non-empty. So this follows immediately from Proposition 7.9.
Proof of Theorem 1.13. From Corollary 7.10, $H^{2}(M, g)=0$, so the actual moduli space is locally isomorphic to $H^{1} / H^{0}$. Depending upon the action of $H^{0}$, the moduli space could therefore be of dimension $\operatorname{dim}\left(H^{1}\right), \operatorname{dim}\left(H^{1}\right)-1$, or $\operatorname{dim}\left(H^{1}\right)-2$. The result then follows immediately from the determination of $H^{1}(M, g)$ in Theorem 7.7.

### 7.5. Final remarks

We end with a non-rigorous remark on the number-theoretic condition appearing in Theorem 7.7. Figure 7.1 contains a plot of the function

$$
\begin{equation*}
H(r, q, p(j))=\left\{\frac{p}{q r}\right\}-\left\{\frac{q^{-1 ; r} p}{r}\right\} \tag{7.51}
\end{equation*}
$$

for $r=3$ and $q=7$, where the horizontal axis indexes the $j$ th prime. The plot begins at the fifth prime, 11, and ends with the 100th prime 541. This, along with other empirical examples, indicates that the cases $H>0$ and $H<0$ occur with approximately the same frequency.


Fig. 7.1. $H(3,7, p(j))$

Acknowledgments. The authors would like to thank John Lott for crucial discussions about index theory. Thanks are also due to Nobuhiro Honda for many valuable discussions on the moduli space of anti-self-dual metrics. Finally, the authors would like to thank the anonymous referee for a very careful reading of the manuscript, and for suggesting many improvements in the exposition.

Research partially supported by NSF Grants DMS-0804042 and DMS-1105187.

## References

[AV12] Ache, A. G., Viaclovsky, J. A.: Obstruction-flat asymptotically locally Euclidean metrics. Geom. Funct. Anal. 22, 832-877 (2012) Zbl 1255.53027 MR 2984119
[Apo90] Apostol, T. M.: Modular Functions and Dirichlet Series in Number Theory. 2 ed., Grad. Texts in Math. 41, Springer, New York (1990) Zbl 0697.10023 MR 1027834
[AG02] Apostolov, V., Gauduchon, P.: Selfdual Einstein Hermitian four-manifolds. Ann. Scuola Norm. Sup. Pisa Cl. Sci. (5) 1, 203-243 (2002) Zbl 1072.53006 MR 1994808
[Aro57] Aronszajn, N.: A unique continuation theorem for solutions of elliptic partial differential equations or inequalities of second order. J. Math. Pures Appl. (9) 36, 235-249 (1957) Zbl 0084.30402 MR 0092067
[AI08] Ashikaga, T., Ishizaka, M.: Another form of the reciprocity law of Dedekind sum. Hokkaido Univ. EPrints Server, no. 908, http://eprints3.math.sci.hokudai.ac.jp/1849/ (2008)
[Bry01] Bryant, R. L.: Bochner-Kähler metrics. J. Amer. Math. Soc. 14, 623-715 (2001) Zbl 1006.53019 MR 1824987
[CS04] Calderbank, D. M. J., Singer, M. A.: Einstein metrics and complex singularities. Invent. Math. 156, 405-443 (2004) Zbl 1061.53026 MR 2052611
[CLW08] Chen, X., Lebrun, C., Weber, B.: On conformally Kähler, Einstein manifolds. J. Amer. Math. Soc. 21, 1137-1168 (2008) Zbl 1208.53072 MR 2425183
[DG06] David, L., Gauduchon, P.: The Bochner-flat geometry of weighted projective spaces. In: Perspectives in Riemannian Geometry, CRM Proc. Lecture Notes 40, Amer. Math. Soc., Providence, RI, 109-156 (2006) Zbl 1109.32019 MR 2237108
[Der83] Derdziński, A.: Self-dual Kähler manifolds and Einstein manifolds of dimension four. Compos. Math. 49, 405-433 (1983) Zbl 0527.53030 MR 0707181
[Dim92] Dimca, A.: Singularities and Topology of Hypersurfaces. Universitext, Springer, New York (1992) Zbl 0753.57001 MR 1194180
[Don89] Donnelly, H.: Rapid decay of eigenfunctions. In: Differential Geometry (Peñíscola, 1988), Lecture Notes in Math. 1410, Springer, Berlin, 161-176 (1989) Zbl 0699.58019 MR 1034278
[EH79] Eguchi, T., Hanson, A. J.: Self-dual solutions to Euclidean gravity. Ann. Phys. 120, 82106 (1979) Zbl 0409.53020 MR 0540896
[GH78] Gibbons, G. W., Hawking, S. W.: Gravitational multi-instantons. Phys. Lett. B 78, 430432 (1978)
[HZ74] Hirzebruch, F., Zagier, D.: The Atiyah-Singer Theorem and Elementary Number Theory. Math. Lecture Ser. 3, Publish or Perish, Boston, MA (1974) Zbl 0288.10001 MR 0650832
[Hit79] Hitchin, N. J.: Polygons and gravitons. Math. Proc. Cambridge Philos. Soc. 85, 465-476 (1979) Zbl 0405.53016 MR 0520463
[Hit97] Hitchin, N. J.: Einstein metrics and the eta-invariant. Boll. Un. Mat. Ital. B (7) 11, no. 2, suppl., 95-105 (1997) Zbl 0973.53519 MR 1456253
[Hon13] Honda, N.: Deformation of LeBrun's ALE metrics with negative mass. Comm. Math. Phys. 322, 127-148 (2013) Zbl 1275.53044 MR 3073160
[Ito93] Itoh, M.: Moduli of half conformally flat structures. Math. Ann. 296, 687-708 (1993) Zbl 0788.58011 MR 1233492
[Ito95] Itoh, M.: The Weitzenböck formula for the Bach operator. Nagoya Math. J. 137, 149181 (1995) Zbl 0819.53017 MR 1324547
[Joy95] Joyce, D. D.: Explicit construction of self-dual 4-manifolds. Duke Math. J. 77, 519-552 (1995) Zbl 0855.57028 MR 1324633
[Kat87] Katase, K.: On the value of Dedekind sums and eta-invariants for 3-dimensional lens spaces. Tokyo J. Math. 10, 327-347 (1987) Zbl 0636.10020 MR 0926246
[Kaw81] Kawasaki, T.: The index of elliptic operators over $V$-manifolds. Nagoya Math. J. 84, 135-157 (1981) Zbl 0437.58020 MR 0641150
[KK92] King, A. D., Kotschick, D.: The deformation theory of anti-self-dual conformal structures. Math. Ann. 294, 591-609 (1992) Zbl 0765.58005 MR 1190446
[Kro89] Kronheimer, P. B.: A Torelli-type theorem for gravitational instantons. J. Differential Geom. 29, 685-697 (1989) Zbl 0671.53046 MR 0992335
[LeB88] LeBrun, C.: Counter-examples to the generalized positive action conjecture. Comm. Math. Phys. 118, 591-596 (1988) Zbl 0659.53050 MR 0962489
[LM08] LeBrun, C., Maskit, B.: On optimal 4-dimensional metrics. J. Geom. Anal. 18, 537-564 (2008) Zbl 1154.53027 MR 2393270
[McC02] McCullough, D.: Isometries of elliptic 3-manifolds. J. London Math. Soc. (2) 65, 167182 (2002) Zbl 1012.57023 MR 1875143
[Rad54] Rademacher, H.: Generalization of the reciprocity formula for Dedekind sums. Duke Math. J. 21, 391-397 (1954) Zbl 0057.03801 MR 0062765
[RG72] Rademacher, H., Grosswald, E.: Dedekind Sums. Carus Math. Monogr. 16, Math. Assoc. Amer., Washington, DC (1972) Zbl 0251.10020 MR 0357299
[Sch84] Schoen, R.: Conformal deformation of a Riemannian metric to constant scalar curvature. J. Differential Geom. 20, 479-495 (1984) Zbl 0576.53028 MR 0788292
[Str10] Streets, J.: Asymptotic curvature decay and removal of singularities of Bach-flat metrics. Trans. Amer. Math. Soc. 362, 1301-1324 (2010) Zbl 1220.53050 MR 2563730
[TV05] Tian, G., Viaclovsky, J.: Moduli spaces of critical Riemannian metrics in dimension four. Adv. Math. 196, 346-372 (2005) Zbl 1252.53045 MR 2166311
[Via13] Viaclovsky, J. A.: An index theorem on anti-self-dual orbifolds. Int. Math. Res. Notices 2013, 3911-3930 Zbl 06438738 MR 3096914
[Wri11] Wright, D.: Compact anti-self-dual orbifolds with torus actions. Selecta Math. (N.S.) 17, 223-280 (2011) Zbl 1245.53045 MR 2803843
[Zag72] Zagier, D. B.: Equivariant Pontrjagin Classes and Applications to Orbit Spaces. Applications of the $G$-Signature Theorem to Transformation Groups, Symmetric Products and Number Theory. Lecture Notes in Math. 290, Springer, Berlin (1972) Zbl 0238.57013 MR 0339202


[^0]:    M. T. Lock: Department of Mathematics, The University of Texas at Austin, Austin, TX 78712, USA; e-mail: mlock@math.utexas.edu
    J. A. Viaclovsky: Department of Mathematics, University of Wisconsin, Madison, WI 53706, USA; e-mail: jeffv@math.wisc.edu

